S-ZDES: ZONAL DETACHED EDDY SIMULATION COUPLED WITH STEADY RANS IN THE WALL REGION

Lars Davidson, www.tfd.chalmers.se/~lada

▲ロト ▲団ト ▲ヨト ▲ヨト 三目 - のへで

DES — DETACHED-EDDY SIMULATIONS

- Problem:
 - the flow in the RANS region is highly unsteady (i.e. URANS)
 - this means that RANS turbulence models (developed for steady flow) are not accurate



DES — DETACHED-EDDY SIMULATIONS

- Problem:
 - the flow in the RANS region is highly unsteady (i.e. URANS)
 - this means that RANS turbulence models (developed for steady flow) are not accurate
- Solution:
 - solve the steady equations in the RANS region



TWO SOLVERS IN THE ENTIRE DOMAIN



Grey color indicates the solver that drives the flow

www.tfd.chalmers.se/~lada

CHALMERS

Workshop, 29 August 2018 3 / 23

DRIFT TERMS ARE ADDED IN WHITE REGIONS



$$\langle \phi(t)
angle_T = rac{1}{T} \int_{-\infty}^t \phi(\tau) \exp(-(t-\tau)/T) d au \Rightarrow$$

 $\langle \phi
angle_T^t \equiv \langle \phi
angle_T = a \langle \phi
angle_T^{t-\Delta t} + (1-a) \phi^t$
 $a = \exp(-\Delta t/T).$

CHALMERS

A (10) A (10) A (10)

WHY IS THE RANS SOLVER STEADY?



- The RANS solver is called every 10th timestep (can probably be called less often)
- The solution in the RANS solver stays steady when the drift term, S_i^{RANS} is steady (constant in time)
- If the integration time T is too small, there will slightly different steady RANS flow every 10th timestep
 - Solution: make the steady RANS solver unsteady but use the large timestep, i.e. 10∆t_{DES}

CHALMERS

PREVIOUS WORK

- The present method is similar to those in [1, 2, 3]. The main differences are that
 - In [1, 3] they use one additional drift terms in the LES momentum equations to control resolved Reynolds stresses
 - ► They include drift terms also in the k and ε equations [1] or the k equation [3].
 - In [1, 3] they include five tuning constants in all drift terms. I have one (*T*).

4 3 5 4 3 5 5

TURBULENCE MODELS

Steady RANS solver



 EARSM (Explicit Algebraic Stress Model) [4] coupled to Wilcox k – ω model [5]



- DES $k \omega$ model
- Lengthscale in dissipation term of the *k* eq.is taken from the IDDES model [6, 7]

RANS and DES turbulence models

NUMERICAL METHOD: CALC-LES & CALC-BFC

CALC-LES [8]: DES solver

- Incompressible finite volume method
- Pressure-velocity coupling treated with fractional step
- Central differencing scheme for momentum eqns
- Hybrid 1st order upwind/2nd order central scheme $k \& \omega$ eqns.
- 2nd-order Crank-Nicholson for time discretization
- CALC-BFC [9]: RANS solver, called every 10th timestep
 - Incompressible finite volume method
 - SIMPLEC
 - MUSCL: 2nd order bounded upwind scheme for momentum eqns
 - Hybrid 1st order upwind/2nd order central scheme $k \& \omega$ eqns.

通 と く ヨ と く ヨ と

FIRST TEST CASE: CHANNEL FLOW

- Reynolds number is $Re_{\tau} = 5200$.
- A 32 \times 96 \times 32 mesh is used
- $x_{max} = 3.2$, $z_{max} = 1.6$, 15% stretching in y direction



CHANNEL FLOW: VELOCITY



 $T = 50\delta/U_b$ — : DES; – – : RANS; •: DNS. Vertical black lines show DES interface.

CHALMERS

THE 1 A

CHANNEL FLOW: TURBULENT QUANTITIES



CHALMERS

SECOND TEST CASE: HUMP FLOW



The domain of the hump. $z_{max} = 0.2$.

- The Reynolds number of the hump flow is $Re_c = 936000$.
- The mesh has $386 \times 120 \times 32$ cells (x, y, z)
- Grid from NASA workshop.¹
- Inlet is located at x = -2.1 and the outlet at x = 4.0,

www.tfd.chalmers.se/~lada

CHALMERS

HUMP FLOW: NUMERICAL ISSUES

Steady RANS solver



- The drift term in the RANS solver in the LES region (white region) causes unphysical oscillations in the skinfriction
- The problem was traced to the source term in the pressure correction equation, the continuity error \dot{m}
- Hence, \dot{m} was set to zero in the LES region.
- As a consequence, the RANS velocity field is driven by the drift term, but the RANS pressure is not correct in this region

→ ∃ → < ∃ →</p>

HUMP FLOW: $C_p \& C_f$



 $T = 20h/U_{in}$. ----: S-DES, $j_0 = 33$; ----: S-DES, $j_0 = 53$; ----: DES

CHALMERS

HUMP FLOW: VELOCITIES



-----: S-ZDES, *j*₀ = 33; ---: S-ZDES, *j*₀ = 53; ----: DES; •: exp [11, 12]; +,+: DES interface.

www.tfd.chalmers.se/~lada

CHALMERS

TURBULENT VISCOSITY (EARSM). S-ZDES, $j_0 = 53$



www.tfd.chalmers.se/~lada

CHALMERS

Workshop, 29 August 2018 16 / 23

HUMP FLOW: SHEAR STRESSES. S-ZDES, $j_0 = 53$



- : DES solver, resolved; - - : RANS solver; - - : DES solver, modeled.

www.tfd.chalmers.se/~lada

CHALMERS

Workshop, 29 August 2018 17 / 23

Hump flow: wrong p in RANS region



www.tfd.chalmers.se/~lada

CHALMERS

Workshop, 29 August 2018 18 / 23

CONCLUSIONS

- A new steady RANS coupled to DES (S-ZDES) is proposed.
- Very good results ...
- but the hump results are maybe/probablby contaminated by a numerical fix
- Drawback: it is dependent on the lower limit of integration time, T for the hump flow
 - $T = 10h/U_{in}$ too small (*h* is hump height)
 - T = 20 and 50 give indentical results
 - For T = 100 we must more than double developing+sampling time to 345 + 345 (7.3 + 7.3 throughflow times)

- ロ ト - (理 ト - (三 ト - (三 ト -)

REFERENCES I

[1] H. Xiao and P. Jenna.

A consistent dual-mesh framework for hybrid LES/RANS modeling.

Journal of Computational Physics, 231:1848–1865, 2012.

[2] B. de Laage de Meux, B. Audebert, R. Manceau, and R. Perrin. Anisotropic linear forcing for synthetic turbulence generation in large eddy simulation and hybrid RANS/LES modeling. *Physics of Fluids A*, 27(035115), 2015.

[3] R. Tunstall, D. Laurence, R. Prosser, and A. Skillen. Towards a generalised dual-mesh hybrid les/rans framework with improved consistency. *Computers & Fluids*, 157:73–83, 2017.

A (10) × (10) × (10) ×

REFERENCES II

 S. Wallin and A. V. Johansson.
 A new explicit algebraic Reynolds stress model for incompressible and compressible turbulent flows.
 Journal of Fluid Mechanics, 403:89–132, 2000.

[5] D. C. Wilcox.

Reassessment of the scale-determining equation. *AIAA Journal*, 26(11):1299–1310, 1988.

[6] M. L. Shur, P. R. Spalart, M. Kh. Strelets, and A. K. Travin. A hybrid RANS-LES approach with delayed-DES and wall-modelled LES capabilities.

International Journal of Heat and Fluid Flow, 29:1638–1649, 2008.

A (10) A (10)

REFERENCES III

 S. Arvidson, L. Davidson, and S.-H. Peng. Interface methods for grey-area mitigation in turbulence-resolving hybrid RANS-LES.

International Journal of Heat and Fluid Flow, 73:236–257, 2018.

[8] L. Davidson.

CALC-LES: A Fortran code for LES and hybrid LES-RANS. Technical report, Division of Fluid Dynamics, Dept. of Mechanics and Maritime Sciences, Chalmers University of Technology, Gothenburg, 2018.

[9] L. Davidson and B. Farhanieh.

CALC-BFC: A finite-volume code employing collocated variable arrangement and cartesian velocity components for computation of fluid flow and heat transfer in complex three-dimensional geometries.

REFERENCES IV

Rept. 95/11, Dept. of Thermo and Fluid Dynamics, Chalmers University of Technology, Gothenburg, 1995.

[10] M. Lee and R. D. Moser.

Direct numerical simulation of turbulent channel flow up to $Re_{\tau} \approx 5200.$

Journal of Fluid Mechanics, 774:395–415, 2015.

[11] D. Greenblatt, K. B. Paschal, C.-S. Yao, J. Harris, N. W. Schaeffler, and A. E. Washburn.

A separation control CFD validation test case. Part 1: Baseline & steady suction.

AIAA-2004-2220, 2004.

[12] D. Greenblatt, K. B. Paschal, C.-S. Yao, and J. Harris. A separation control CFD validation test case Part 1: Zero efflux oscillatory blowing. AIAA-2005-0485, 2005.

www.tfd.chalmers.se/~lada

CHALMERS

Workshop, 29 August 2018 23 / 23

3