## L.1. Part I: Data of Two-dimensional flow

What about the viscous terms: where do they play an important role? Which terms are non-zero *at* the wall? (you can easily show that with paper and pen).

So far we have looked at the  $\bar{v}_1$ -momentum equation. The database corresponds to a two-dimensional flow. Now let's think of the forces as vectors. The gradient of the normal stresses in the  $x_1 - x_2$  plane represent the force vector (see Eq. L.1)

$$\mathbf{F}_{N} = \left(-\frac{\partial \overline{v_{1}^{\prime 2}}}{\partial x_{1}}, -\frac{\partial \overline{v_{2}^{\prime 2}}}{\partial x_{2}}\right) \tag{L.2}$$

and the corresponding force vector due to the shear stresses reads (see Eq. L.1)

$$\mathbf{F}_{S} = \left(-\frac{\partial \overline{v_{1}'v_{2}'}}{\partial x_{2}}, -\frac{\partial \overline{v_{1}'v_{2}'}}{\partial x_{1}}\right)$$
(L.3)

Find the first term in Eqs. L.2 and L.3 in the  $\bar{v}_1$  momentum equation, Eq. L.1. Consider the second line in Eq. L.1 which is the  $\bar{v}_2$  equation and find the other two terms in Eqs. L.2 and L.3. Note that  $\mathbf{F}_N$  and  $\mathbf{F}_S$  are forces per unit volume  $([N/m^3])$ .

Assignment 1.3. The left-hand side can be formulated in three different ways (steady flow),  $d\bar{v}_i/dt$ ,  $\bar{v}_j\partial\bar{v}_i/\partial x_j$  and  $\partial/\partial x_j(\bar{v}_i\bar{v}_j)$ , see Section 2.4. Consider the first form,  $d\bar{v}_i/dt$ . This is the Lagrangian form, i.e. we follow a particle, see Fig. 1.1. Now, pick a vertical grid line at  $x_1$  and try to estimate  $d\bar{v}_1/dt$ . Then you want to estimate  $v_1$ ,  $\Delta t$  seconds later. During this time the fluid particle has moved the vector distance  $(\Delta x_1, \Delta x_2)$  to the next vertical grid line with the speed  $V = (v_1^2 + v_2^2)^{1/2}$  so that  $\Delta t = \Delta s/V$  where  $\Delta s = ((\Delta x_1)^2 + (\Delta x_2)^2)^{1/2}$ 

Compare this way of computing the left-hand side with  $\partial/\partial x_i(\bar{v}_1\bar{v}_i)$ .

Note that we have now estimated  $d\bar{v}_1/dt$  using the discrete approximation of the chain rule in Eq. 1.1, i.e.

$$\frac{d\bar{v}_1}{dt} \simeq \frac{\Delta\bar{v}_1}{\Delta t} = V \frac{\Delta\bar{v}_1}{\Delta s} \simeq \bar{v}_1 \frac{\partial\bar{v}_1}{\partial x_1} + \bar{v}_2 \frac{\partial\bar{v}_1}{\partial x_2} \tag{L.4}$$

- Assignment 1.4. Plot the vector field  $\mathbf{F}_N$  to learn something about its properties. When  $v_2'^2$  reaches a maximum or a minimum along a grid line normal to the wall, what happens with the vector field  $\mathbf{F}_N$ ? Zoom-in on interesting regions.
- Assignment 1.5. Plot also vector fields of the shear stress,  $\mathbf{F}_S$  (see Eq. L.3), the pressure gradient and the viscous terms. Zoom up in interesting regions. Anything interesting?

## L.1.3 The turbulent kinetic energy equation

The exact transport equation for the turbulent kinetic energy, k, reads

$$\frac{\partial}{\partial x_j} (\bar{v}_j k) = \nu \frac{\partial^2 k}{\partial x_j \partial x_j} + P^k + D_k - \varepsilon$$

$$P^k = -\overline{v'_i v'_j} \frac{\partial \bar{v}_i}{\partial x_j}$$
(L.5)