2013-08-26, Exam in

Turbulence modeling, MTF270

- Time: 8.30–12.30 Location: M
- Teacher: Lars Davidson, phone 772 1404, 0730-791 161
- Aids
 - Formula sheet, 2 pages, appended to the exam
- Checking the evaluation and results of your written exam: contact the assistant, Hamid Abedi
- Grading: 20-29p: 3, 30-39: 4, 40-50: 5.
- T1. a) Derive the Boussinesq assumption. (5p)
 - b) Discuss and show how the dissipation term, ε_{ij} , is modeled. (5p)
- T2. a) Describe the physical effect of stable stratification and unstable stratification on turbulence. (5p)
 - b) Consider streamline curvature for a streamline formed as a circular arc (convex curvature). Show that the turbulence is dampened if $\partial v_{\theta}/\partial r > 0$ and that it is enhanced if the sign of $\partial v_{\theta}/\partial r$ is negative. (5p)
- T3. a) Derive a transport equation for ω from the k and ε transport equations. (5p) Hint:

$$D_{\omega}^{T} \simeq \frac{2\nu_{t}}{\sigma_{\varepsilon}k} \frac{\partial k}{\partial x_{i}} \frac{\partial \omega}{\partial x_{j}} + \frac{\partial}{\partial x_{j}} \left(\frac{\nu_{t}}{\sigma_{\varepsilon}} \frac{\partial \omega}{\partial x_{j}} \right)$$

b) Consider the energy spectrum. Show the three different regions (the large energy-containing scales, the -5/3 range and the dissipating scales). Where should the cut-off be located? Show where the SGS scales, grid (i.e resolved) scales and the cut-off, κ_c are located in the spectrum.

- T4. a) Show that when a first-order upwind schemes is used for the convection term, an additional diffusion term and dissipation terms appear because of a numerical SGS viscosity.
 - b) The modified (reduced) length scale in two-equation DES models can be introduced in different equations. Which equations and which term? What is the effect on the modeled, turbulent quantities?
- T5. a) Describe URANS. How is the instantaneous velocity decomposed? What turbulence models are used? What is scale separation?
 - b) Describe the PANS model. What is the main modification compared to the standard $k \varepsilon$ model? What is the physical meaning of f_k ? Describe what happens to the equation system when f_k is reduced. (5p)

MTF270 Turbulence modelling: Formula sheet

August 21, 2013

The continuity, Navier-Stokes and temperature equations for incompressible flow with constant viscosity read (*conservative* form)

$$\frac{\partial v_i}{\partial x_i} = 0$$

$$\rho_0 \frac{\partial v_i}{\partial t} + \rho_0 \frac{\partial v_i v_j}{\partial x_j} = -\frac{\partial p}{\partial x_i} + \mu \frac{\partial^2 v_i}{\partial x_j \partial x_j} - \rho_0 \beta (\theta - \theta_0) g_i$$

$$\frac{\partial \theta}{\partial t} + \frac{\partial v_i \theta}{\partial x_i} = \alpha \frac{\partial^2 \theta}{\partial x_i \partial x_i}$$

The Navier-Stokes equation for incompressible flow with constant viscosity read (*non-conservative* form, p denotes the hydrostatic pressure, i.e. p = 0 if $v_i = 0$)

$$\rho_0 \frac{\partial v_i}{\partial t} + \rho_0 v_j \frac{\partial v_i}{\partial x_j} = -\frac{\partial p}{\partial x_i} + \mu \frac{\partial^2 v_i}{\partial x_j \partial x_j}$$

The time averaged continuity equation, Navier-Stokes equation temperature equations read

$$\frac{\partial \bar{v}_i}{\partial x_i} = 0$$

$$\rho_0 \frac{\partial \bar{v}_i \bar{v}_j}{\partial x_j} = -\frac{\partial \bar{p}}{\partial x_i} + \frac{\partial}{\partial x_j} \left(\mu \frac{\partial \bar{v}_i}{\partial x_j} - \rho_0 \overline{v_i' v_j'} \right) - \rho_0 \beta (\bar{\theta} - \theta_0) g_i$$

$$\frac{\partial \bar{v}_i \bar{\theta}}{\partial x_i} = \alpha \frac{\partial^2 \bar{\theta}}{\partial x_i \partial x_i} - \frac{\partial \overline{v_i' \theta'}}{\partial x_i}$$

The Boussinesq assumption reads

$$\overline{v_i'v_j'} = -\nu_t \left(\frac{\partial \overline{v}_i}{\partial x_j} + \frac{\partial \overline{v}_j}{\partial x_i} \right) + \frac{2}{3} \delta_{ij} k = -2\nu_t \overline{s}_{ij} + \frac{2}{3} \delta_{ij} k$$

The modeled $\overline{v_i'v_j'}$ equation with IP model reads

$$\bar{v}_{k} \frac{\partial \overline{v'_{i}v'_{j}}}{\partial x_{k}} = \text{ (convection)}$$

$$-\overline{v'_{i}v'_{k}} \frac{\partial \bar{v}_{j}}{\partial x_{k}} - \overline{v'_{j}v'_{k}} \frac{\partial \bar{v}_{i}}{\partial x_{k}} \text{ (production)}$$

$$-c_{1} \frac{\varepsilon}{k} \left(\overline{v'_{i}v'_{j}} - \frac{2}{3}\delta_{ij}k \right) \text{ (slow part)}$$

$$-c_{2} \left(P_{ij} - \frac{2}{3}\delta_{ij}P^{k} \right) \text{ (rapid part)}$$

$$+c_{1w}\rho_{0} \frac{\varepsilon}{k} \left[\overline{v'_{k}v'_{m}}n_{k}n_{m}\delta_{ij} - \frac{3}{2}\overline{v'_{i}v'_{k}}n_{k}n_{j} \right]$$

$$-\frac{3}{2}\overline{v'_{j}v'_{k}}n_{k}n_{i} \right] f \left[\frac{\ell_{t}}{x_{n}} \right] \text{ (wall, slow part)}$$

$$+c_{2w} \left[\Phi_{km,2}n_{k}n_{m}\delta_{ij} - \frac{3}{2}\Phi_{ik,2}n_{k}n_{j} \right]$$

$$-\frac{3}{2}\Phi_{jk,2}n_{k}n_{i} \right] f \left[\frac{\ell_{t}}{x_{n}} \right] \text{ (wall, rapid part)}$$

$$+\nu \frac{\partial^{2}\overline{v'_{i}v'_{j}}}{\partial x_{k}\partial x_{k}} \text{ (viscous diffusion)}$$

$$+\frac{\partial}{\partial x_{k}} \left[c_{k} \overline{v'_{k}v'_{m}} \frac{k}{\varepsilon} \frac{\partial \overline{v'_{i}v'_{j}}}{\partial x_{m}} \right] \text{ (turbulent diffusion)}$$

$$-g_{i}\beta \overline{v'_{j}\theta'} - g_{j}\beta \overline{v'_{i}\theta'} \text{ (buoyancy production)}$$

$$-\frac{2}{3}\varepsilon\delta_{ij} \text{ (dissipation)}$$

▶The exact transport equation for turbulent heat heat flux vector $\overline{v_i'\theta'}$ reads

$$\frac{\partial \overline{v_i'\theta'}}{\partial t} + \frac{\partial}{\partial x_k} \overline{v_k} \overline{v_i'\theta'} = -\overline{v_i'v_k'} \frac{\partial \overline{\theta}}{\partial x_k} - \overline{v_k'\theta'} \frac{\partial \overline{v_i}}{\partial x_k} - \overline{\frac{\theta'}{\rho} \frac{\partial p'}{\partial x_i}} - \frac{\partial}{\partial x_k} \overline{v_k'v_i'\theta'}$$

$$+ \alpha \overline{\frac{\partial}{\partial x_k} \left(v_i' \frac{\partial \theta'}{\partial x_k} \right)} + \nu \overline{\frac{\partial}{\partial x_k} \left(\theta' \frac{\partial v_i'}{\partial x_k} \right)} - \underline{\left(\nu + \alpha \right) \frac{\partial v_i'}{\partial x_k} \frac{\partial \theta'}{\partial x_k}} - \underline{g_i \beta \overline{\theta'^2}}_{G_{i\theta}}$$

$$= \frac{\partial \overline{v_i'\theta'}}{\partial x_k} \overline{v_k'v_i'\theta'} - \underline{v_k'\theta'} \overline{v_k'\theta'} \overline{v_k'\theta'} \overline{v_k'\theta'} \overline{v_k'\phi'} \overline{$$

▶ The exact k equation reads

$$\frac{\partial k}{\partial t} + \frac{\partial \bar{v}_j k}{\partial x_j} = -\overline{v_i' v_j'} \frac{\partial \bar{v}_i}{\partial x_j} - \frac{\partial}{\partial x_j} \left[\frac{1}{\rho} \overline{v_j' p'} + \frac{1}{2} \overline{v_j' v_i' v_i'} - \nu \frac{\partial k}{\partial x_j} \right] - \nu \overline{\frac{\partial v_i'}{\partial x_j} \frac{\partial v_i'}{\partial x_j}}$$

▶The exact $\overline{v_i'v_j'}$ equation reads

$$\frac{\partial \overline{v_i'v_j'}}{\partial t} + \frac{\partial}{\partial x_k} (\bar{v}_k \overline{v_i'v_j'}) = -\overline{v_j'v_k'} \frac{\partial \bar{v}_i}{\partial x_k} - \overline{v_i'v_k'} \frac{\partial \bar{v}_j}{\partial x_k} - \frac{\partial}{\partial v_i'v_j'} \frac{\partial \bar{v}_j}{\partial x_k} - \frac{\partial}{\partial v_i'v_j'} \frac{\partial}{\partial v_i'v_j'} + \frac{1}{\rho} \delta_{ik} \overline{v_j'p'} - \nu \frac{\partial \overline{v_i'v_j'}}{\partial x_k} + \frac{1}{\rho} \overline{v_j'} \frac{\partial v_j'}{\partial x_j} - 2\nu \frac{\partial v_j'}{\partial x_k} \frac{\partial v_j'}{\partial x_k} \frac{\partial v_j'}{\partial x_k} + \frac{\partial v_j'}{\partial x_k} \frac{\partial v_j'}{\partial x_k} + \frac{\partial v_j'}{\partial x_k} \frac{\partial v_j'}{\partial x_k} \frac{\partial v_j'}{\partial x_k} + \frac{\partial v_j'}{\partial x_k} \frac{\partial v_j'}{\partial x_k} \frac{\partial v_j'}{\partial x_k} + \frac{\partial v_j'}{\partial x_k} \frac{\partial v_j'}{\partial x_k} \frac{\partial v_j'}{\partial x_k} + \frac{\partial v_j'}{\partial x_k} + \frac{\partial v_j'}{\partial x_k} \frac{\partial$$

▶ The modeled k and ε equations

$$\begin{split} \frac{\partial k}{\partial t} + \bar{v}_j \frac{\partial k}{\partial x_j} &= \nu_t \left(\frac{\partial \bar{v}_i}{\partial x_j} + \frac{\partial \bar{v}_j}{\partial x_i} \right) \frac{\partial \bar{v}_i}{\partial x_j} + g_i \beta \frac{\nu_t}{\sigma_\theta} \frac{\partial \bar{\theta}}{\partial x_i} \\ &- \varepsilon + \frac{\partial}{\partial x_j} \left[\left(\nu + \frac{\nu_t}{\sigma_k} \right) \frac{\partial k}{\partial x_j} \right] \\ \frac{\partial \varepsilon}{\partial t} + \bar{v}_j \frac{\partial \varepsilon}{\partial x_j} &= \frac{\varepsilon}{k} c_{\varepsilon 1} \nu_t \left(\frac{\partial \bar{v}_i}{\partial x_j} + \frac{\partial \bar{v}_j}{\partial x_i} \right) \frac{\partial \bar{v}_i}{\partial x_j} \\ &+ c_{\varepsilon 1} g_i \frac{\varepsilon}{k} \frac{\nu_t}{\sigma_\theta} \frac{\partial \bar{\theta}}{\partial x_i} - c_{\varepsilon 2} \frac{\varepsilon^2}{k} + \frac{\partial}{\partial x_j} \left[\left(\nu + \frac{\nu_t}{\sigma_\varepsilon} \right) \frac{\partial \varepsilon}{\partial x_j} \right] \end{split}$$