

Electric welding arc modeling
with the 3D solver OpenFOAM
- **A comparison of different electromagnetic models** -

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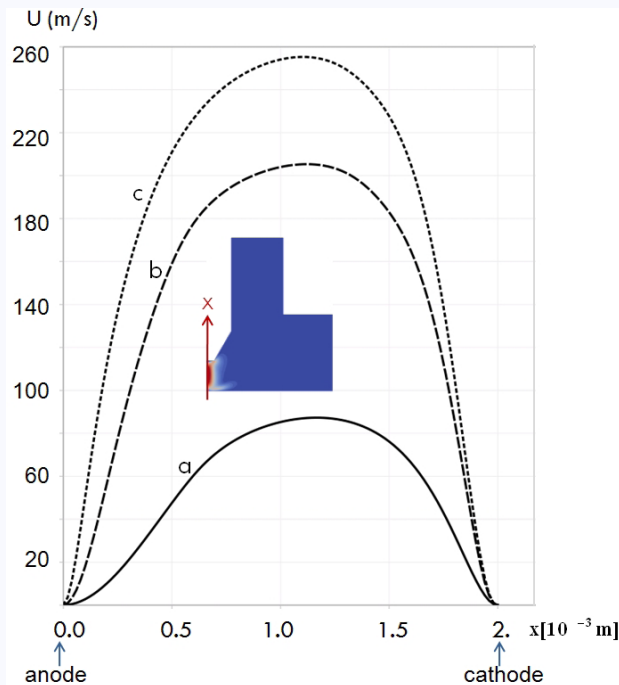
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Gothenburg, Sweden

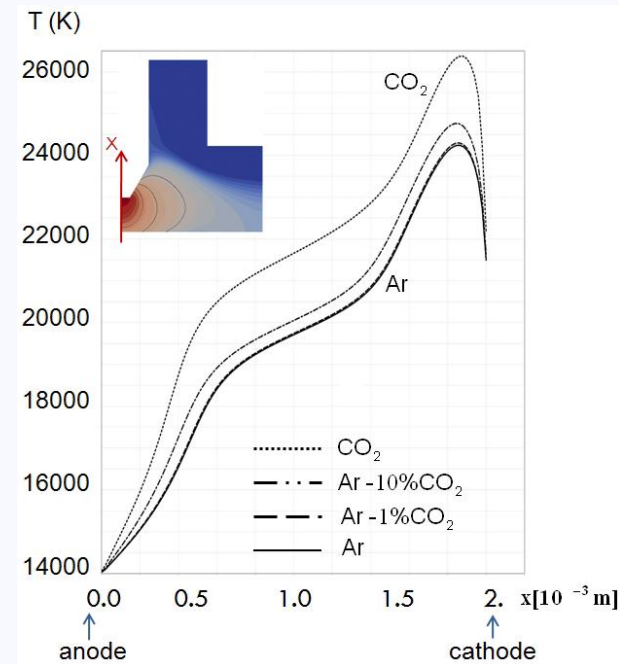
- Context / Motivation:
better understand the heat source
- Software OpenFOAM-1.6.x
 - open source CFD software
 - C++ library of object-oriented classes
for implementing solvers for continuum mechanics

IIW intermediate meeting, Trollhättan, Doc. XII-2017-11

” Numerical simulation of $\text{Ar}-x\%\text{CO}_2$ shielding gas and its effect on an electric welding arc”



Influence of BC on anode and cathode composition



Influence of gas

Here focus on a comparison of different electromagnetic models

Model: thermal fluid part

Main assumptions (plasma core):

- one-fluid model
- local thermal equilibrium
- mechanically incompressible and thermally expansible
- *steady flow*
- *laminar flow (assuming laminar shielding gas inlet)*

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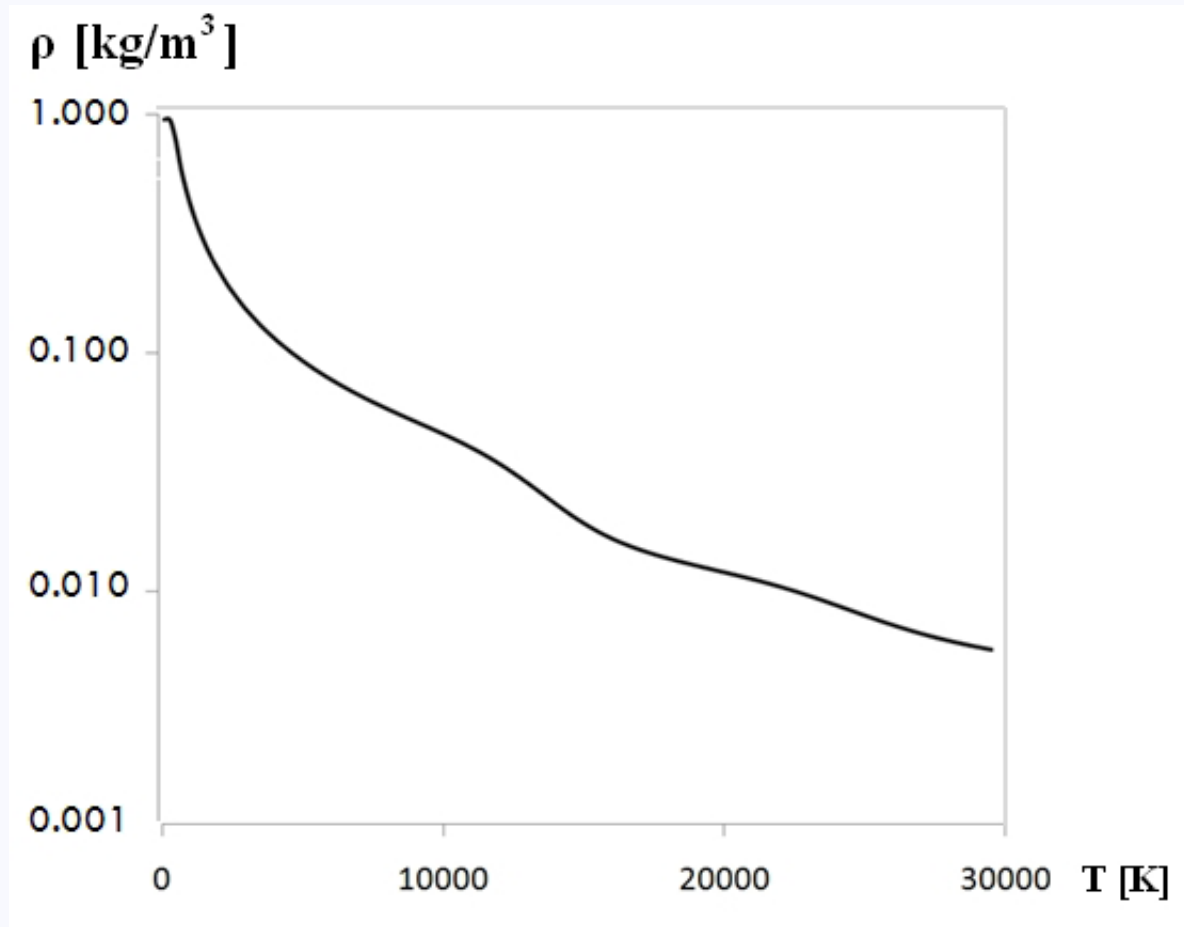
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Argon plasma density as function of temperature.

Model: thermal fluid part

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- plasma optically thin

- *steady*

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Model: thermal fluid part

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- one-fluid model
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Model: electromagnetic part

- 3D model with
 - Electric field formulation
 - Magnetic field formulation
- 2D axi-symmetric models:
 - ✓ Electric potential formulation
 - ✓ Magnetic field formulation

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- 3D model with $V, \vec{A} \rightarrow \vec{E}, \vec{J}, \vec{B}$
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- 3D model with $V, \vec{A} \Rightarrow \vec{E}, \vec{J}, \vec{B}$
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Model: electromagnetic part

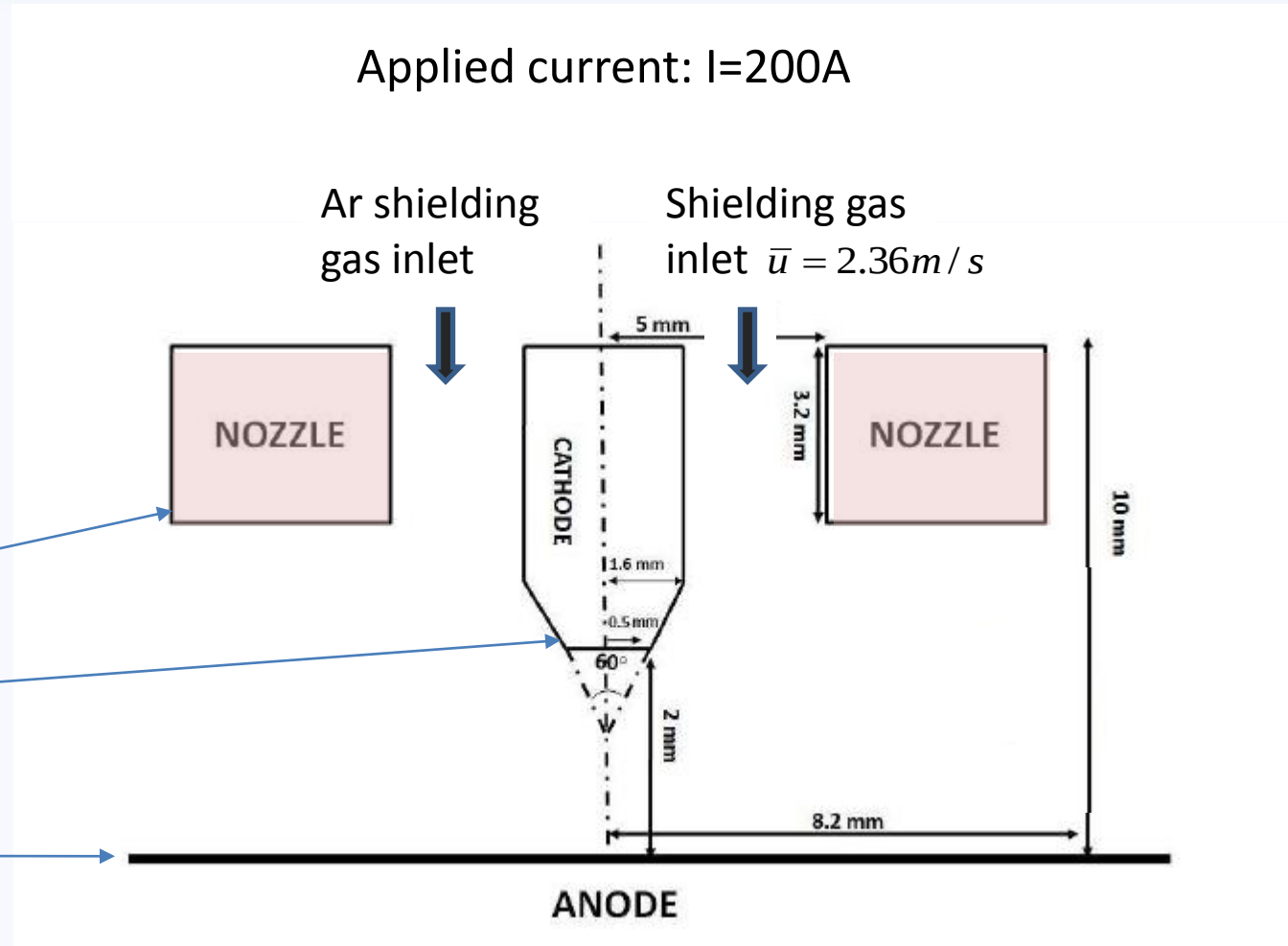
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M.A. Ramírez , G. Trapaga and J. McKelliget (2003). A comparison between two different numerical formulations of welding arc simulation. *Modelling Simul. Mater. Sci. Eng*, **11**, pp. 675-695.

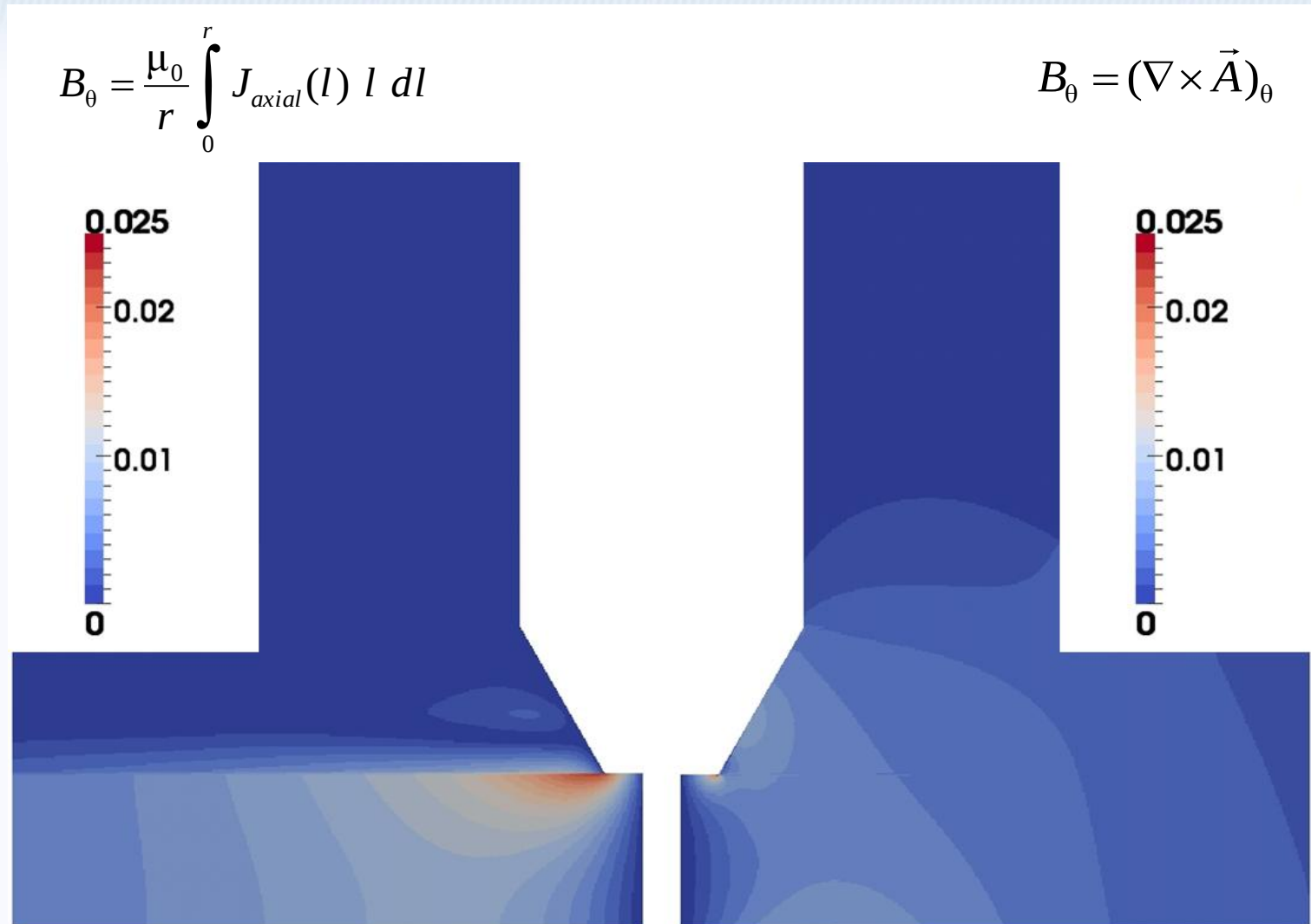
Test case: Tungsten Inert Gas welding



Picture of a TIG torch



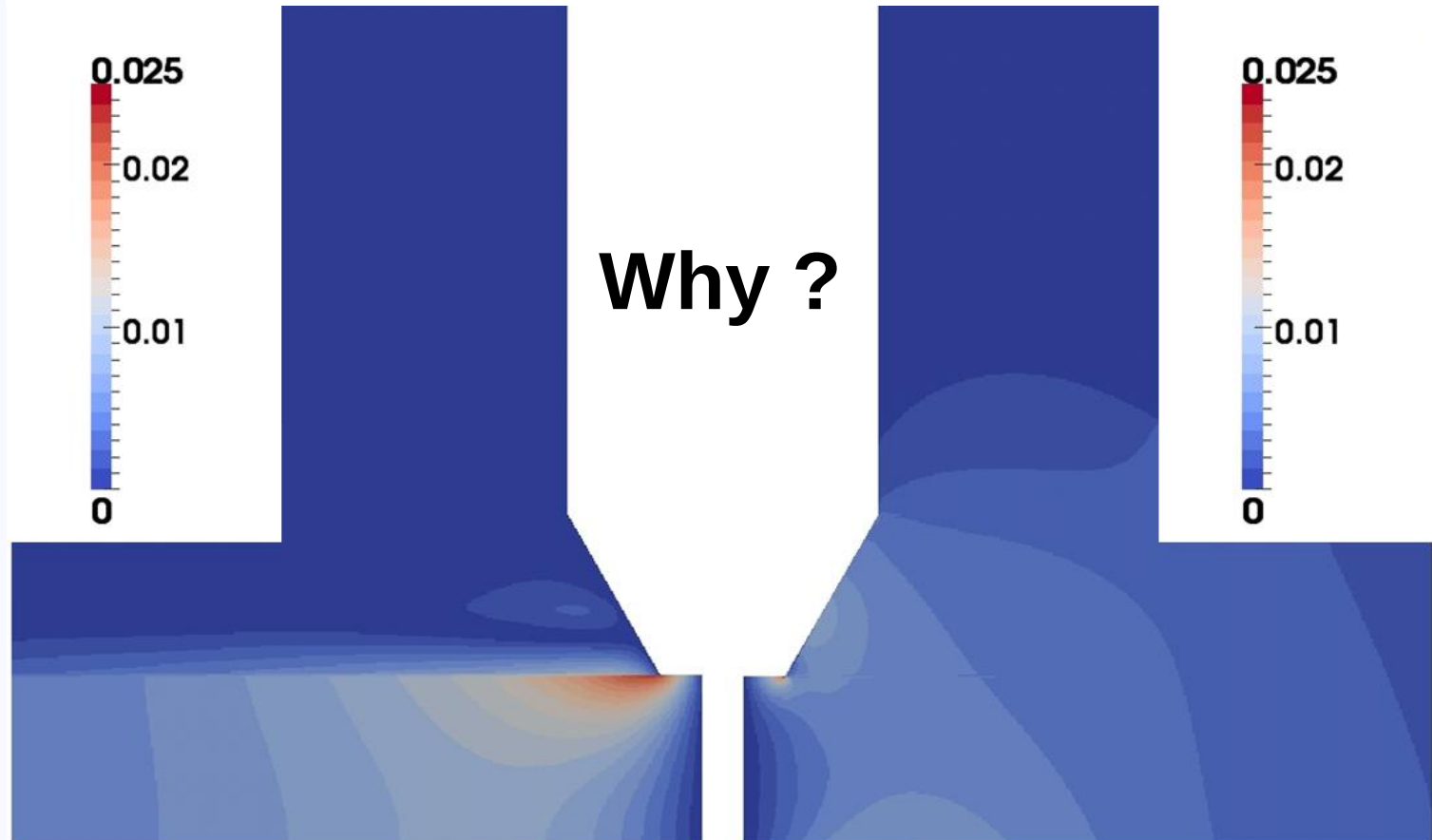
Sketch of the cross section of a TIG torch



Magnetic field magnitude calculated with the electric potential formulation (left) and the 3D

$$B_{\theta} = \frac{\mu_0}{r} \int_0^r J_{axial}(l) l dl$$

$$B_{\theta} = (\nabla \times \vec{A})_{\theta}$$



Magnetic field magnitude calculated with the electric potential formulation (left) and the 3D

Maxwell's equations

- Gauss' law for magnetism: $\nabla \cdot \vec{B} = 0$

- Ampère's law: $\epsilon_0 \mu_0 \partial_t \vec{E} = \nabla \times \vec{B} - \mu_0 \vec{J}$

- Faraday's law: $\partial_t \vec{B} = -\nabla \times \vec{E}$

- Gauss' law: $\nabla \cdot \vec{E} = \epsilon_0^{-1} q_{tot}$

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- + generalized Ohm's law $\vec{J} = \vec{J}_{drift} + \vec{J}_{ind} + \vec{J}_{Hall} + \vec{J}_{diff} + \vec{J}_{ther}$

Assumptions (plasma core)

- $\lambda_{Debye} \approx 10^{-8} m \ll L_c \Rightarrow$ local electro-neutrality
- quasi-steady electromagnetic phenomena
-
-

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- $\beta_{Lar} = v_{Lar} / v_{coll} \ll 1 \Rightarrow \vec{J}_{Hall} = \frac{-\sigma}{n_e} \vec{J} \times \vec{B} \ll \vec{J}_{drift} = \sigma \vec{E}$

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- $Re_m \ll 1 \Rightarrow \vec{J}_{ind} = \sigma \vec{u} \times \vec{B} \ll \vec{J}_{drift}$
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Electromagnetic model for arc plasma core

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(1) with Lorentz gauge : $\nabla \cdot \vec{A} = 0$

Electromagnetic model for arc plasma core

(1)

$$\nabla \cdot (\sigma \nabla V) = 0$$

$$\Delta \vec{A} = \mu_0 \sigma \nabla V$$

with

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$$\vec{B} = \nabla \times \vec{A}$$

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Electromagnetic model for arc plasma core

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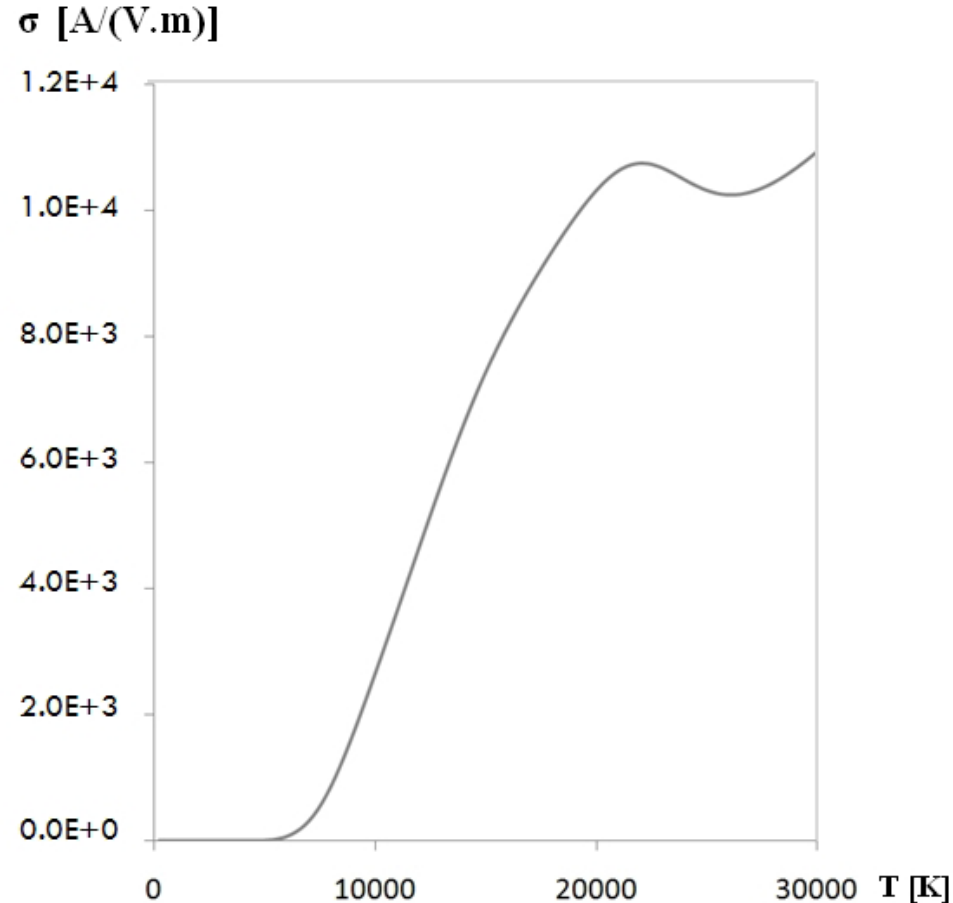
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Argon plasma electric conductivity
as function of temperature

(1) with Lorentz gauge : $\nabla \cdot \vec{A} = 0$

2D axi-symmetric case

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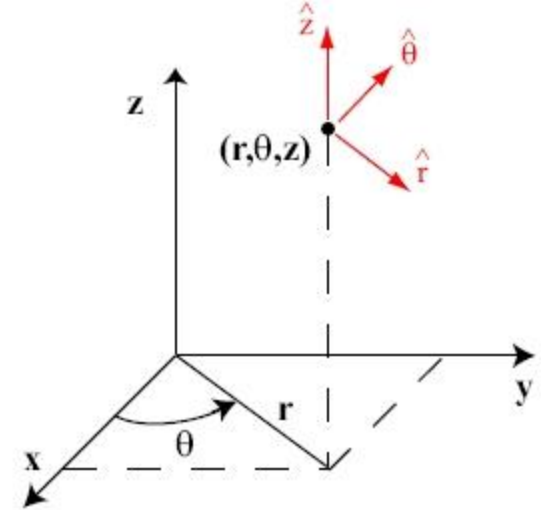
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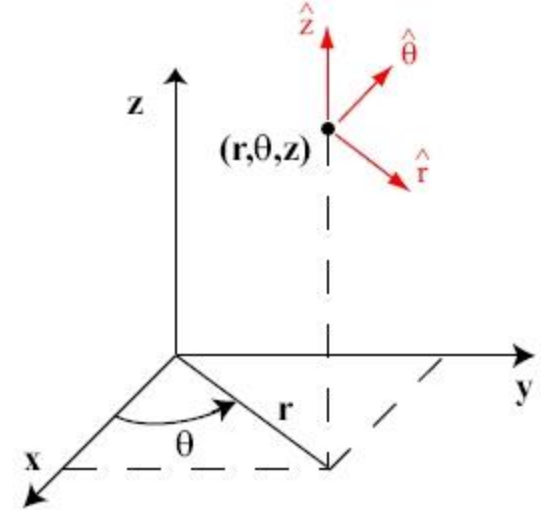
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2D axi-symmetric case⁽¹⁾



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Then



with

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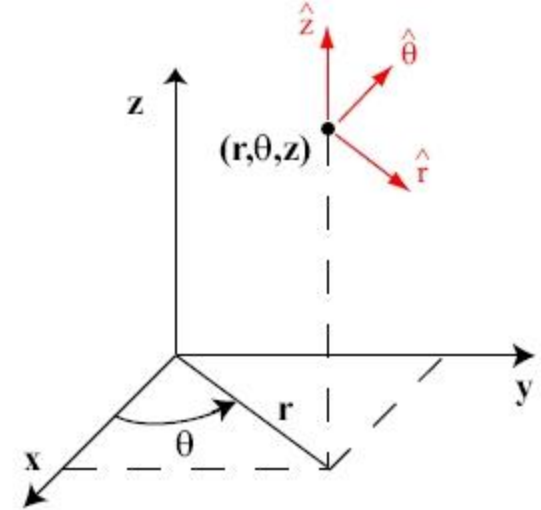
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$$V(r, z)$$



with

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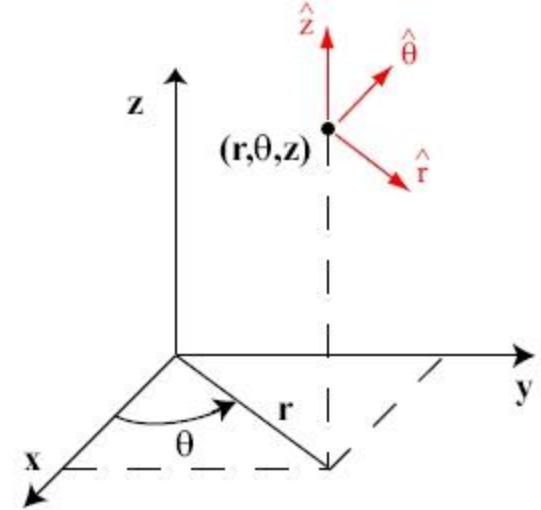


$$V(r, z)$$

$$\vec{A} = (A_r, 0, A_z) \quad \text{with} \quad A_r(r, z), A_z(r, z)$$

with

with



2D axi-symmetric case⁽¹⁾

Then

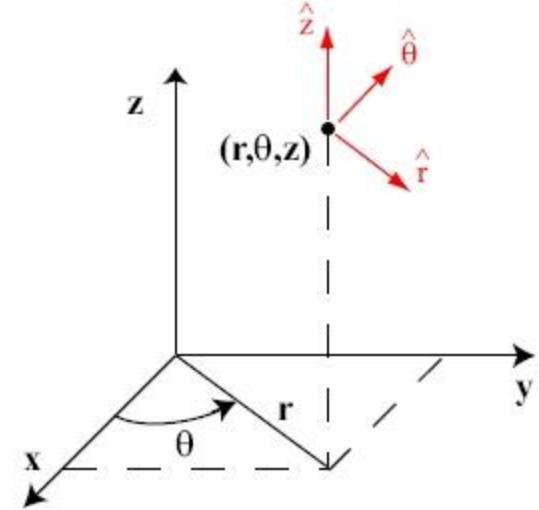


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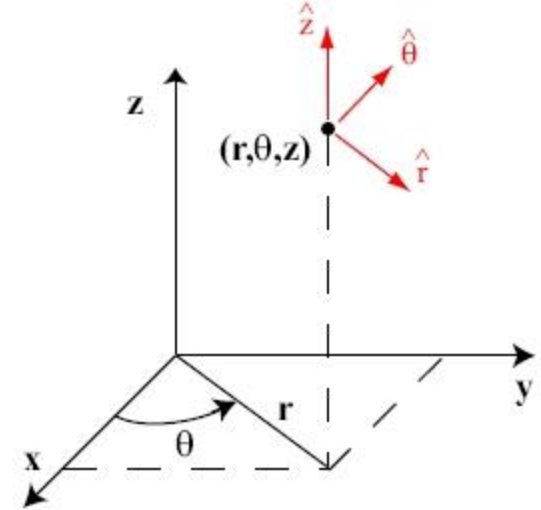
$$\vec{J} = (J_r, 0, J_z) \quad \text{with} \quad J_r(r,z), J_z(r,z)$$

with



2D axi-symmetric case⁽¹⁾

Then



$$V(r, z)$$

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$$\vec{J} = (J_r, 0, J_z) \quad \text{with} \quad J_r(r, z), \quad J_z(r, z)$$

$$\vec{B} = (0, B_\theta, 0) \quad \text{with} \quad B_\theta(r, z)$$

2D axi-symmetric case ⁽¹⁾

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \sigma \frac{\partial V}{\partial r} \right) + \frac{\partial}{\partial z} \left(\sigma \frac{\partial V}{\partial z} \right) = 0$$

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_r}{\partial r} \right) + \frac{\partial^2 A_r}{\partial z^2} - \frac{A_r}{r^2} = \mu_0 \sigma \frac{\partial V}{\partial r}$$

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with

$$\vec{J} : \quad J_r = \sigma E_r = -\sigma \frac{\partial V}{\partial r}$$

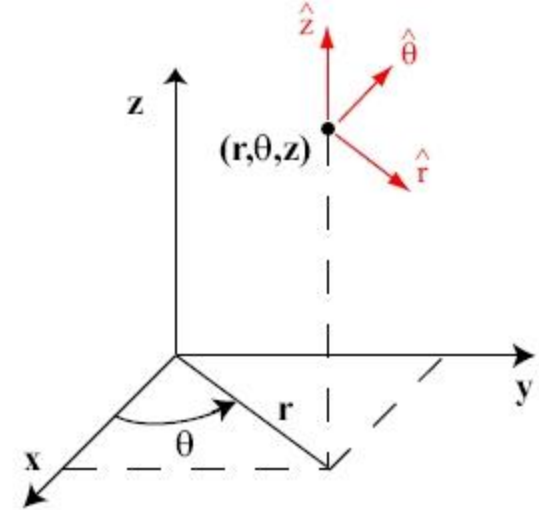
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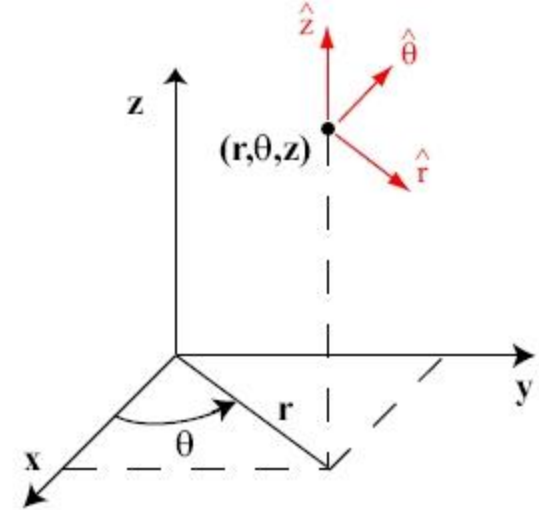
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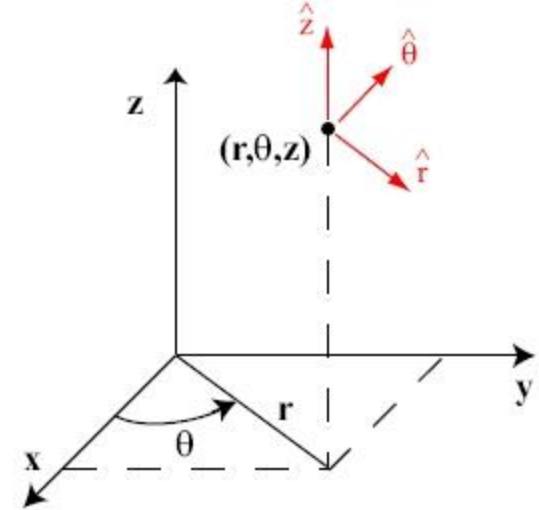
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(1) $\longleftrightarrow$

$$\frac{\partial B_\theta}{\partial z} = \mu_0 \sigma \frac{\partial V}{\partial r} = -\mu_0 J_r$$

$$\frac{1}{r} \frac{\partial (r B_\theta)}{\partial r} = -\mu_0 \sigma \frac{\partial V}{\partial z} = \mu_0 J_z$$

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(1) $\longleftrightarrow$

$$\frac{\partial B_\theta}{\partial z} = \mu_0 \sigma \frac{\partial V}{\partial r} = -\mu_0 J_r$$

$$\frac{1}{r} \frac{\partial(rB_\theta)}{\partial r} = -\mu_0 \sigma \frac{\partial V}{\partial z} = \mu_0 J_z$$

(2) $\searrow$

Induction diffusion equation

$$\frac{\partial}{\partial r} \left(\frac{1}{\sigma r} \frac{\partial(rB_\theta)}{\partial r} \right) + \frac{\partial}{\partial z} \left(\frac{1}{\sigma} \frac{\partial B_\theta}{\partial z} \right) = 0$$

(1) with Lorentz gauge : $\nabla \cdot \vec{A} = 0$

(2) as $\partial_{r,z}^2 V = \partial_{z,r}^2 V$

2D axi-symmetric case – equivalent formulations

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_r}{\partial r} \right) + \frac{\partial^2 A_r}{\partial z^2} - \frac{A_r}{r^2} = \mu_0 \sigma \frac{\partial V}{\partial r}$$

(1) $\longleftrightarrow$

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_z}{\partial r} \right) + \frac{\partial^2 A_z}{\partial z^2} = \mu_0 \sigma \frac{\partial V}{\partial z}$$

$$\frac{\partial B_\theta}{\partial z} = \mu_0 \sigma \frac{\partial V}{\partial r} = -\mu_0 J_r$$

$$\frac{1}{r} \frac{\partial(rB_\theta)}{\partial r} = -\mu_0 \sigma \frac{\partial V}{\partial z} = \mu_0 J_z$$

(2) $\swarrow$

Induction diffusion equation

$$\frac{\partial}{\partial r} \left(\frac{1}{\sigma r} \frac{\partial(rB_\theta)}{\partial r} \right) + \frac{\partial}{\partial z} \left(\frac{1}{\sigma} \frac{\partial B_\theta}{\partial z} \right) = 0$$

(2) $\searrow$

$$B_\theta(r, z) = \int_{l=z_0}^{l=z} J_z(r, l) dl + B_\theta(r, z_0)$$

(1) with Lorentz gauge : $\nabla \cdot \vec{A} = 0$

(2) as $\partial_{r,z}^2 V = \partial_{z,r}^2 V$

2D axi-symmetric case – equivalent formulations

F1

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_r}{\partial r} \right) + \frac{\partial^2 A_r}{\partial z^2} - \frac{A_r}{r^2} = \mu_0 \sigma \frac{\partial V}{\partial r}$$

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_z}{\partial r} \right) + \frac{\partial^2 A_z}{\partial z^2} = \mu_0 \sigma \frac{\partial V}{\partial z}$$

(1)

$$\frac{\partial B_\theta}{\partial z} = \mu_0 \sigma \frac{\partial V}{\partial r} = -\mu_0 J_r$$

$$\frac{1}{r} \frac{\partial(rB_\theta)}{\partial r} = -\mu_0 \sigma \frac{\partial V}{\partial z} = \mu_0 J_z$$

(2)

Induction diffusion equation

F2

$$\frac{\partial}{\partial r} \left(\frac{1}{\sigma r} \frac{\partial(rB_\theta)}{\partial r} \right) + \frac{\partial}{\partial z} \left(\frac{1}{\sigma} \frac{\partial B_\theta}{\partial z} \right) = 0$$

(2)

F3

$$B_\theta(r, z) = \int_{l=z_0}^{l=z} J_z(r, l) dl + B_\theta(r, z_0)$$

(1) with Lorentz gauge : $\nabla \cdot \vec{A} = 0$

(2) as $\partial_{r,z}^2 V = \partial_{z,r}^2 V$

2D axi-symmetric case – equivalent formulations

F1

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_r}{\partial r} \right) + \frac{\partial^2 A_r}{\partial z^2} - \frac{A_r}{r^2} = \mu_0 \sigma \frac{\partial V}{\partial r}$$

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_z}{\partial r} \right) + \frac{\partial^2 A_z}{\partial z^2} = \mu_0 \sigma \frac{\partial V}{\partial z}$$

(1)

$$\frac{\partial B_\theta}{\partial z} = \mu_0 \sigma \frac{\partial V}{\partial r} = -\mu_0 J_r$$

$$\frac{1}{r} \frac{\partial(rB_\theta)}{\partial r} = -\mu_0 \sigma \frac{\partial V}{\partial z} = \mu_0 J_z$$

(2)

F2

Induction diffusion equation

$$\frac{\partial}{\partial r} \left(\frac{1}{\sigma r} \frac{\partial(rB_\theta)}{\partial r} \right) + \frac{\partial}{\partial z} \left(\frac{1}{\sigma} \frac{\partial B_\theta}{\partial z} \right) = 0$$

(2)

F3

$$B_\theta(r, z) = \int_{l=z_0}^{l=z} J_z(r, l) dl + B_\theta(r, z_0)$$

(1) with Lorentz gauge : $\nabla \cdot \vec{A} = 0$

(2) as $\partial_{r,z}^2 V = \partial_{z,r}^2 V$

2D axi-symmetric case – equivalent formulations

F1

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_r}{\partial r} \right) + \frac{\partial^2 A_r}{\partial z^2} - \frac{A_r}{r^2} = \mu_0 \sigma \frac{\partial V}{\partial r}$$

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_z}{\partial r} \right) + \frac{\partial^2 A_z}{\partial z^2} = \mu_0 \sigma \frac{\partial V}{\partial z}$$

(1) $\longleftrightarrow$

$$\frac{\partial B_\theta}{\partial z} = \mu_0 \sigma \frac{\partial V}{\partial r} = -\mu_0 J_r$$

$$\frac{1}{r} \frac{\partial(rB_\theta)}{\partial r} = -\mu_0 \sigma \frac{\partial V}{\partial z} = \mu_0 J_z$$

(2) $\searrow$

F2

Induction diffusion equation

$$\frac{\partial}{\partial r} \left(\frac{1}{\sigma r} \frac{\partial(rB_\theta)}{\partial r} \right) + \frac{\partial}{\partial z} \left(\frac{1}{\sigma} \frac{\partial B_\theta}{\partial z} \right) = 0$$

(2) $\searrow$

F3

$$B_\theta(r, z) = \int_{l=z_0}^{l=z} J_z(r, l) dl + B_\theta(r, z_0)$$

(1) with Lorentz gauge : $\nabla \cdot \vec{A} = 0$

(2) as $\partial_{r,z}^2 V = \partial_{z,r}^2 V$

2D axi-symmetric case – equivalent formulations

F1

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_r}{\partial r} \right) + \frac{\partial^2 A_r}{\partial z^2} - \frac{A_r}{r^2} = \mu_0 \sigma \frac{\partial V}{\partial r}$$

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_z}{\partial r} \right) + \frac{\partial^2 A_z}{\partial z^2} = \mu_0 \sigma \frac{\partial V}{\partial z}$$

(1) $\longleftrightarrow$

$$\frac{\partial B_\theta}{\partial z} = \mu_0 \sigma \frac{\partial V}{\partial r} = -\mu_0 J_r$$

$$\frac{1}{r} \frac{\partial(rB_\theta)}{\partial r} = -\mu_0 \sigma \frac{\partial V}{\partial z} = \mu_0 J_z$$

(2) $\swarrow$

F2

Induction diffusion equation

$$\frac{\partial}{\partial r} \left(\frac{1}{\sigma r} \frac{\partial(rB_\theta)}{\partial r} \right) + \frac{\partial}{\partial z} \left(\frac{1}{\sigma} \frac{\partial B_\theta}{\partial z} \right) = 0$$

(2) $\swarrow$

F3

$$B_\theta(r, z) = \int_{l=z_0}^{l=z} J_z(r, l) dl + B_\theta(r, z_0)$$

(1) with Lorentz gauge : $\nabla \cdot \vec{A} = 0$

(2) as $\partial_{r,z}^2 V = \partial_{z,r}^2 V$

Magnetic field formulation

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_r}{\partial r} \right) + \frac{\partial^2 A_r}{\partial z^2} - \frac{A_r}{r^2} = \mu_0 \sigma \frac{\partial V}{\partial r}$$

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_z}{\partial r} \right) + \frac{\partial^2 A_z}{\partial z^2} = \mu_0 \sigma \frac{\partial V}{\partial z}$$

(1) $\longleftrightarrow$

$$\frac{\partial B_\theta}{\partial z} = \mu_0 \sigma \frac{\partial V}{\partial r} = -\mu_0 J_r$$

$$\frac{1}{r} \frac{\partial(rB_\theta)}{\partial r} = -\mu_0 \sigma \frac{\partial V}{\partial z} = \mu_0 J_z$$

(2) $\searrow$

Induction diffusion equation

$$\frac{\partial}{\partial r} \left(\frac{1}{\sigma r} \frac{\partial(rB_\theta)}{\partial r} \right) + \frac{\partial}{\partial z} \left(\frac{1}{\sigma} \frac{\partial B_\theta}{\partial z} \right) = 0$$

(2) $\swarrow$

$$B_\theta(r, z) = \int_{l=z_0}^{l=z} J_z(r, l) dl + B_\theta(r, z_0)$$

(1) with Lorentz gauge : $\nabla \cdot \vec{A} = 0$

(2) as $\partial_{r,z}^2 V = \partial_{z,r}^2 V$

2D axi-symmetric case ⁽¹⁾

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \sigma \frac{\partial V}{\partial r} \right) + \frac{\partial}{\partial z} \left(\sigma \frac{\partial V}{\partial z} \right) = 0$$

F1

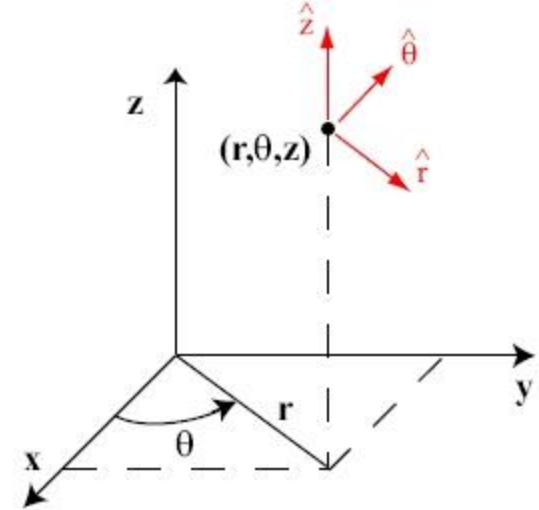
$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_r}{\partial r} \right) + \frac{\partial^2 A_r}{\partial z^2} - \frac{A_r}{r^2} = \mu_0 \sigma \frac{\partial V}{\partial r}$$

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial A_z}{\partial r} \right) + \frac{\partial^2 A_z}{\partial z^2} = \mu_0 \sigma \frac{\partial V}{\partial z}$$

with

$$\vec{J} : \quad J_r = \sigma E_r = -\sigma \frac{\partial V}{\partial r} \quad \text{and} \quad J_z = \sigma E_z = -\sigma \frac{\partial V}{\partial z}$$

$$\vec{B} : \quad B_\theta = \frac{\partial A_r}{\partial z} - \frac{\partial A_z}{\partial r}$$



(1) with Lorentz gauge : $\nabla \cdot \vec{A} = 0$

2D axi-symmetric case⁽¹⁾

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \sigma \frac{\partial V}{\partial r} \right) + \frac{\partial}{\partial z} \left(\sigma \frac{\partial V}{\partial z} \right) = 0$$

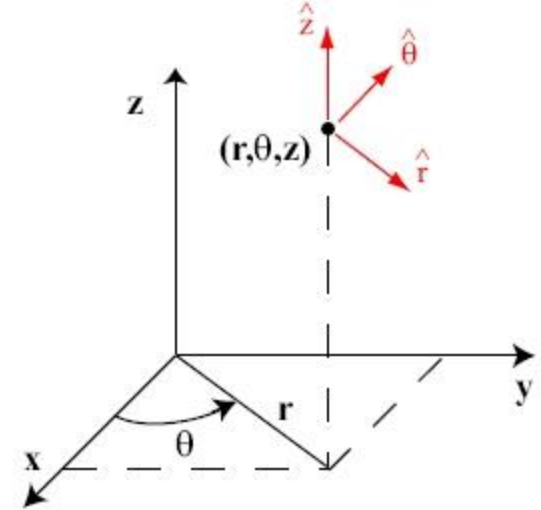
Induction diffusion equation

F2

$$\frac{\partial}{\partial r} \left(\frac{1}{\sigma r} \frac{\partial (r B_\theta)}{\partial r} \right) + \frac{\partial}{\partial z} \left(\frac{1}{\sigma} \frac{\partial B_\theta}{\partial z} \right) = 0$$

with

$$\vec{J} : \quad J_r = \sigma E_r = -\sigma \frac{\partial V}{\partial r} \quad \text{and} \quad J_z = \sigma E_z = -\sigma \frac{\partial V}{\partial z}$$



(1) with Lorentz gauge : $\nabla \cdot \vec{A} = 0$

Electric potential formulation

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \sigma \frac{\partial V}{\partial r} \right) + \frac{\partial}{\partial z} \left(\sigma \frac{\partial V}{\partial z} \right) = 0$$

Induction diffusion equation

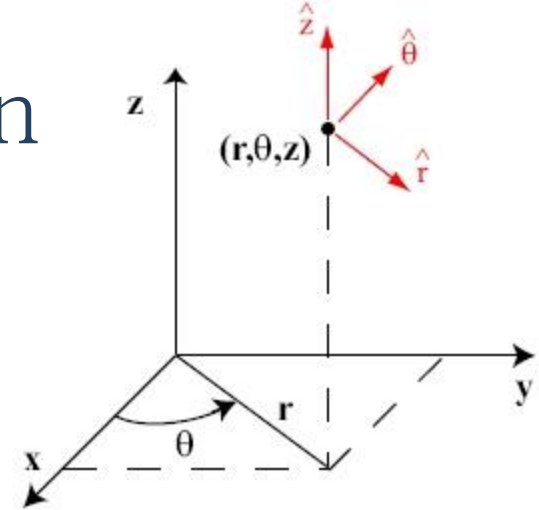
F2

$$\frac{\partial}{\partial r} \left(\frac{1}{\sigma r} \frac{\partial (r B_\theta)}{\partial r} \right) + \frac{\partial}{\partial z} \left(\frac{1}{\sigma} \frac{\partial B_\theta}{\partial z} \right) = 0$$

0

with

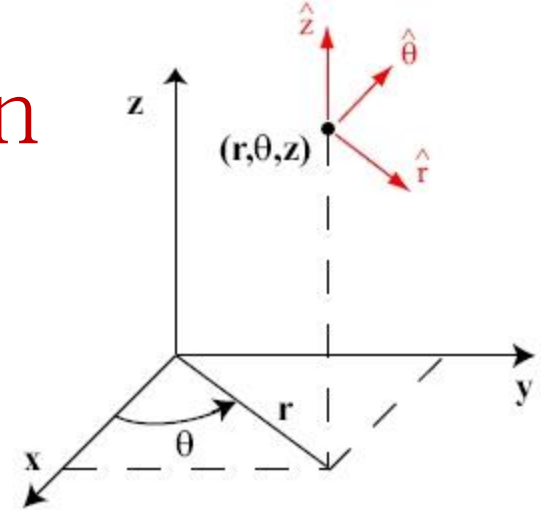
$$\vec{J} : \quad J_r = \sigma E_r = -\sigma \frac{\partial V}{\partial r} \quad \text{and} \quad J_z = \sigma E_z = -\sigma \frac{\partial V}{\partial z}$$



(1) with Lorentz gauge : $\nabla \cdot \vec{A} = 0$

Electric potential formulation

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \sigma \frac{\partial V}{\partial r} \right) + \frac{\partial}{\partial z} \left(\sigma \frac{\partial V}{\partial z} \right) = 0$$



$$B_{\theta}(r) = \frac{\mu_0}{r} \int_{l=r_0}^{l=r} l J_z(l) dl + B_{\theta}(r_0)$$

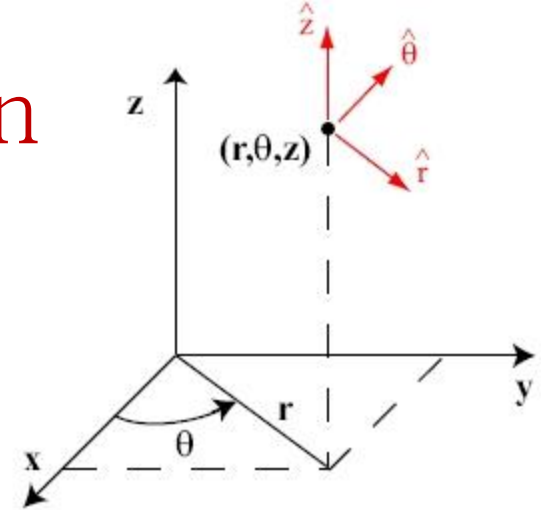
with

$$\vec{J} : \quad J_r = \sigma E_r = -\sigma \frac{\partial V}{\partial r} \quad \text{and} \quad J_z = \sigma E_z = -\sigma \frac{\partial V}{\partial z}$$

(1) with Lorentz gauge : $\nabla \cdot \vec{A} = 0$

Electric potential formulation

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \sigma \frac{\partial V}{\partial r} \right) + \frac{\partial}{\partial z} \left(\sigma \frac{\partial V}{\partial z} \right) = 0$$



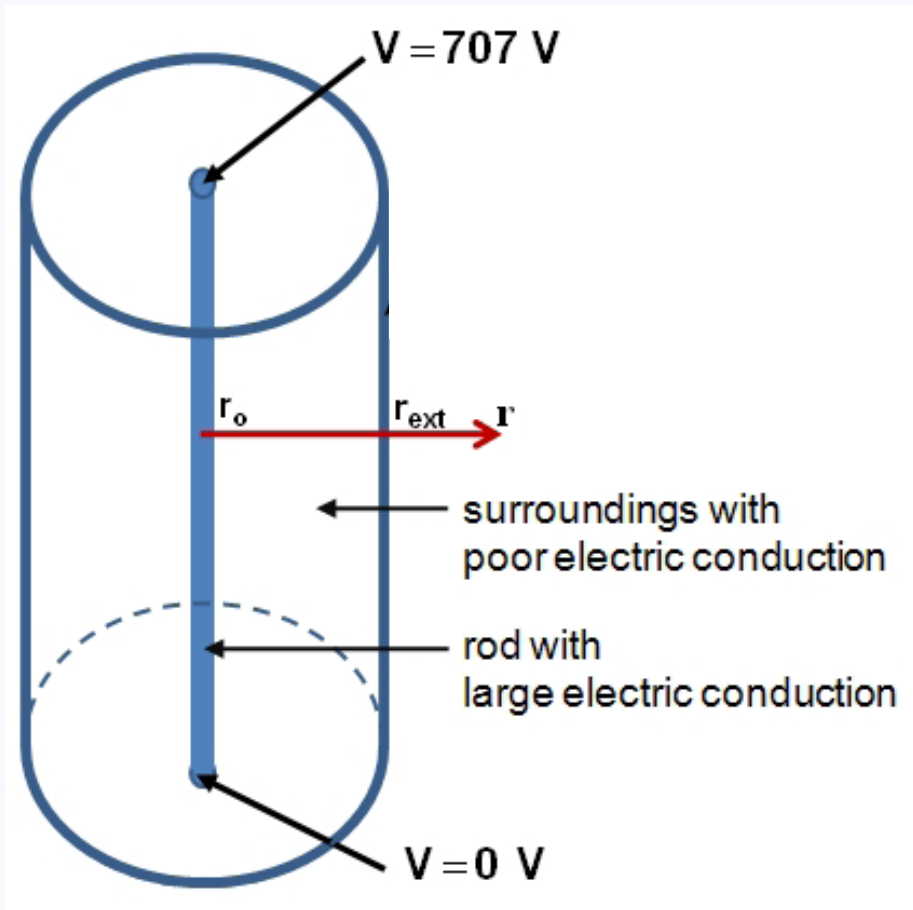
$$B_{\theta}(r) = \frac{\mu_0}{r} \int_{l=r_0}^{l=r} l J_z(l) dl + B_{\theta}(r_0)$$

with

$$\vec{J} : \quad J_r = \sigma E_r = -\sigma \frac{\partial V}{\partial r} \quad \ll \quad J_z = \sigma E_z = -\sigma \frac{\partial V}{\partial z}$$

(1) with Lorentz gauge : $\nabla \cdot \vec{A} = 0$

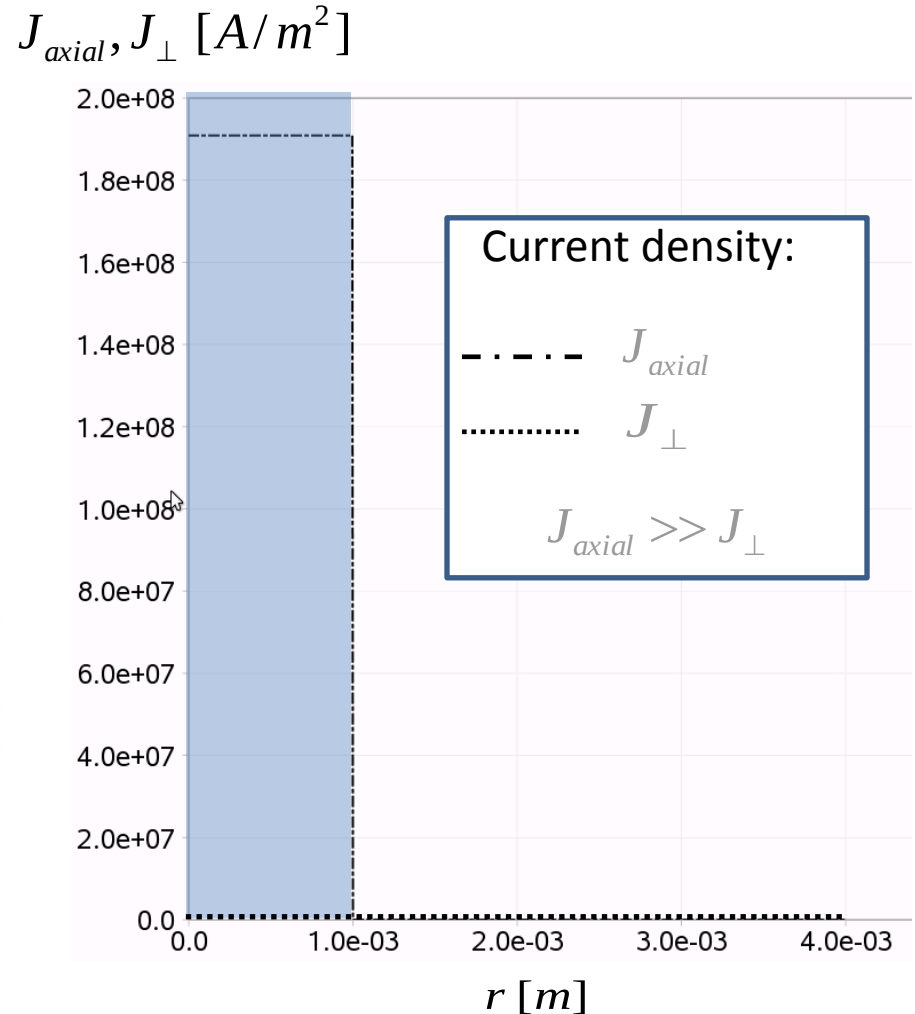
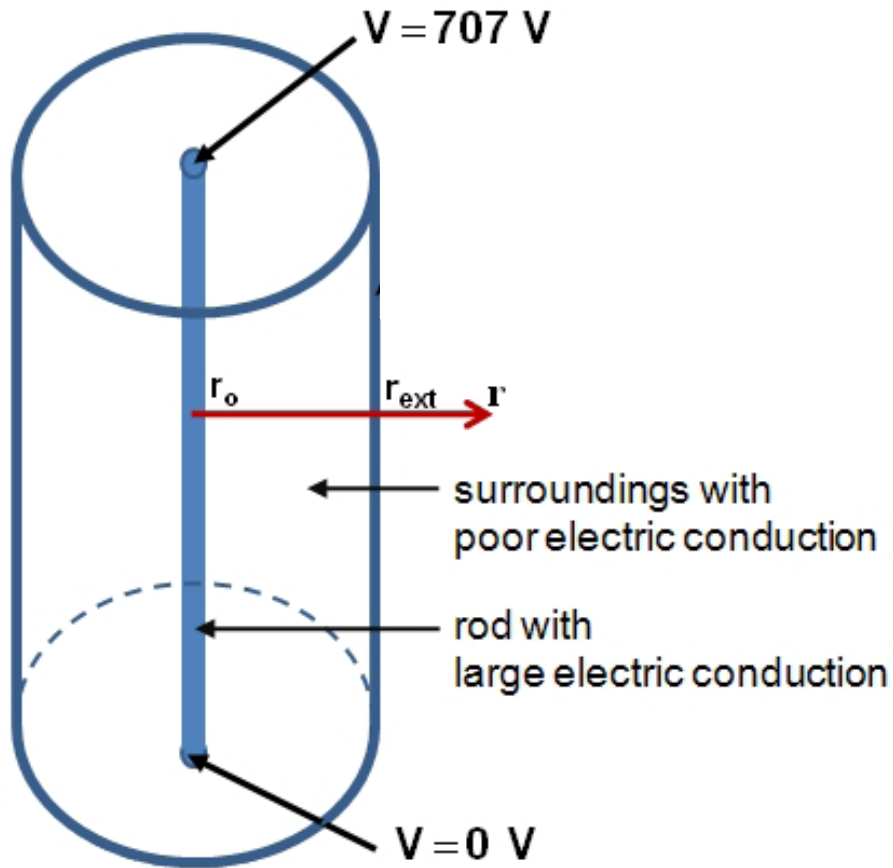
Test case: infinite conducting rod



$\sigma = 2700 \text{ A}/(\text{V.m})$

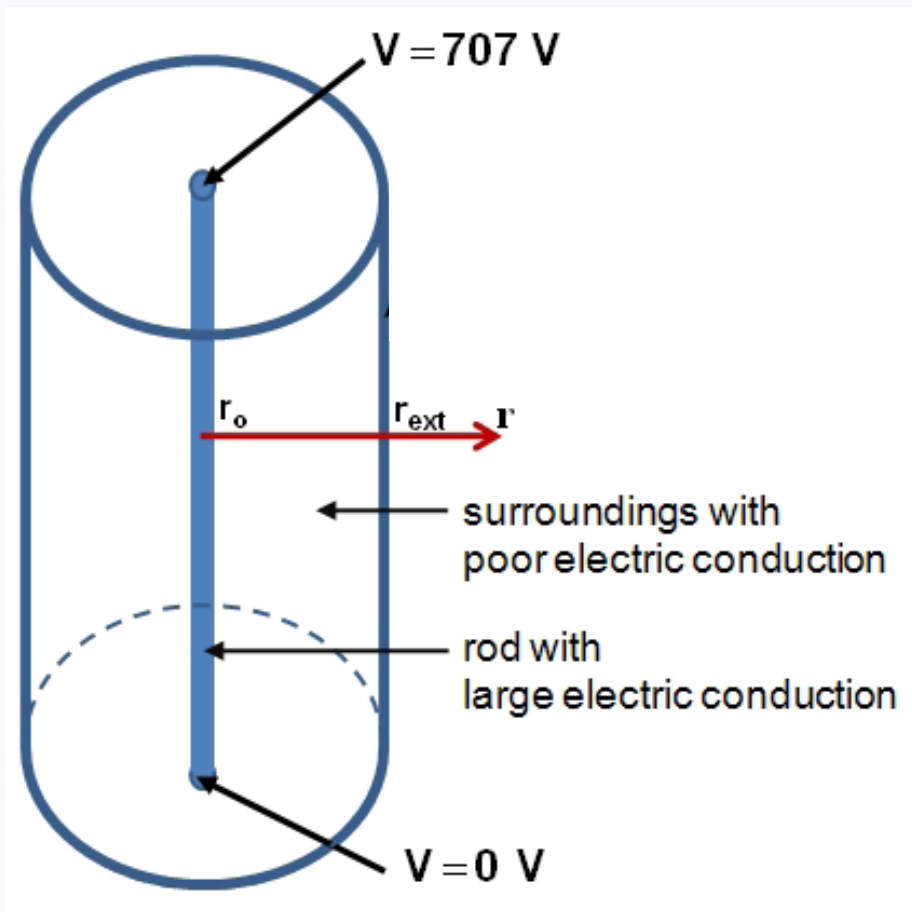
$\sigma = 10^{-5} \text{ A}/(\text{V.m})$

Test case: infinite conducting rod



$\sigma = 2700 \text{ A/(V.m)}$
 $\sigma = 10^{-5} \text{ A/(V.m)}$

Test case: infinite conducting rod

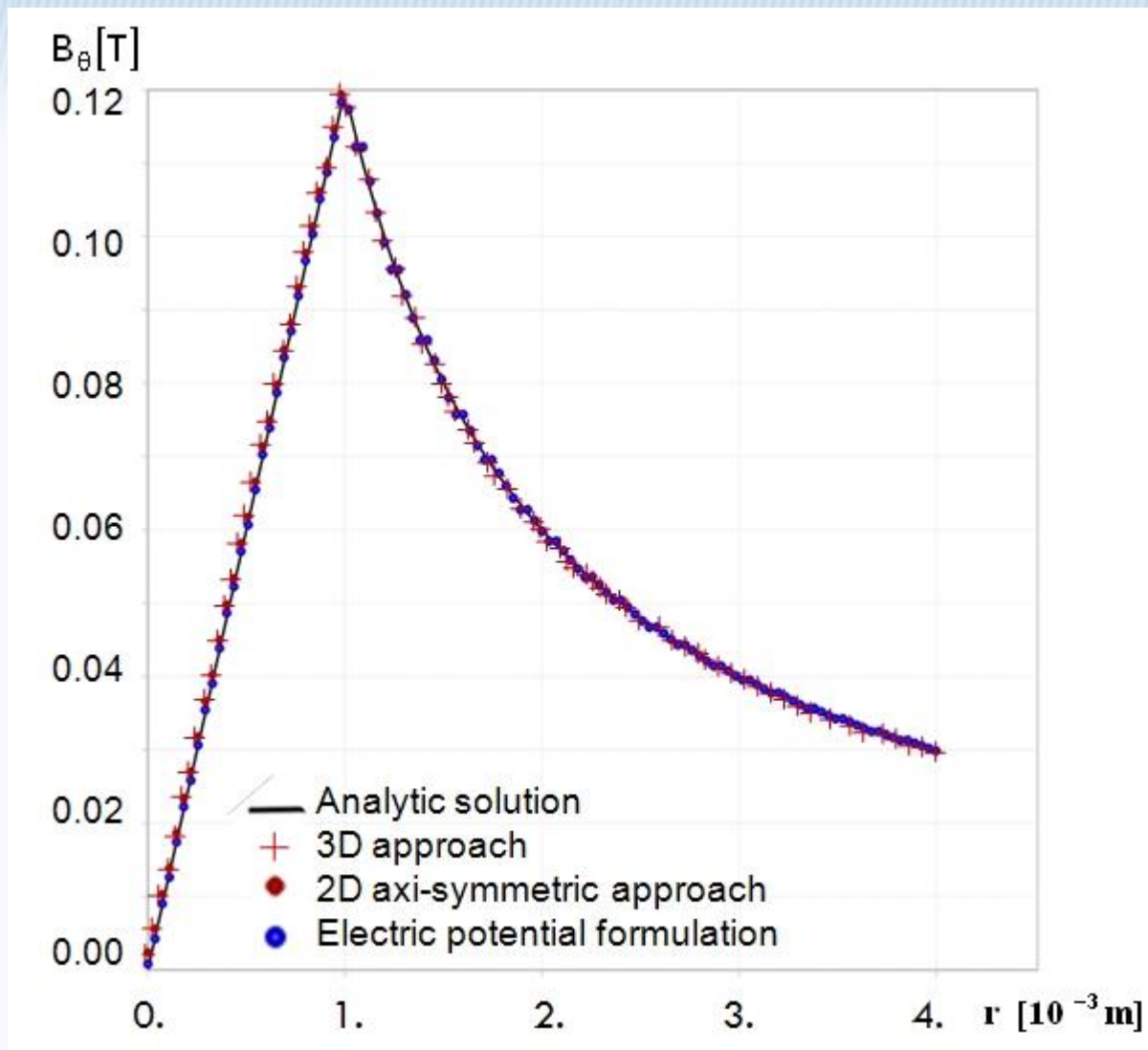


Analytic solution:

$$B_{\theta}(r) = \frac{\mu_o J_{\text{axial}} r}{2} \quad \text{if } r < r_o ,$$

$$B_{\theta}(r) = \frac{\mu_o J_{\text{axial}} r_o^2}{2 r} \quad \text{if } r \geq r_o$$

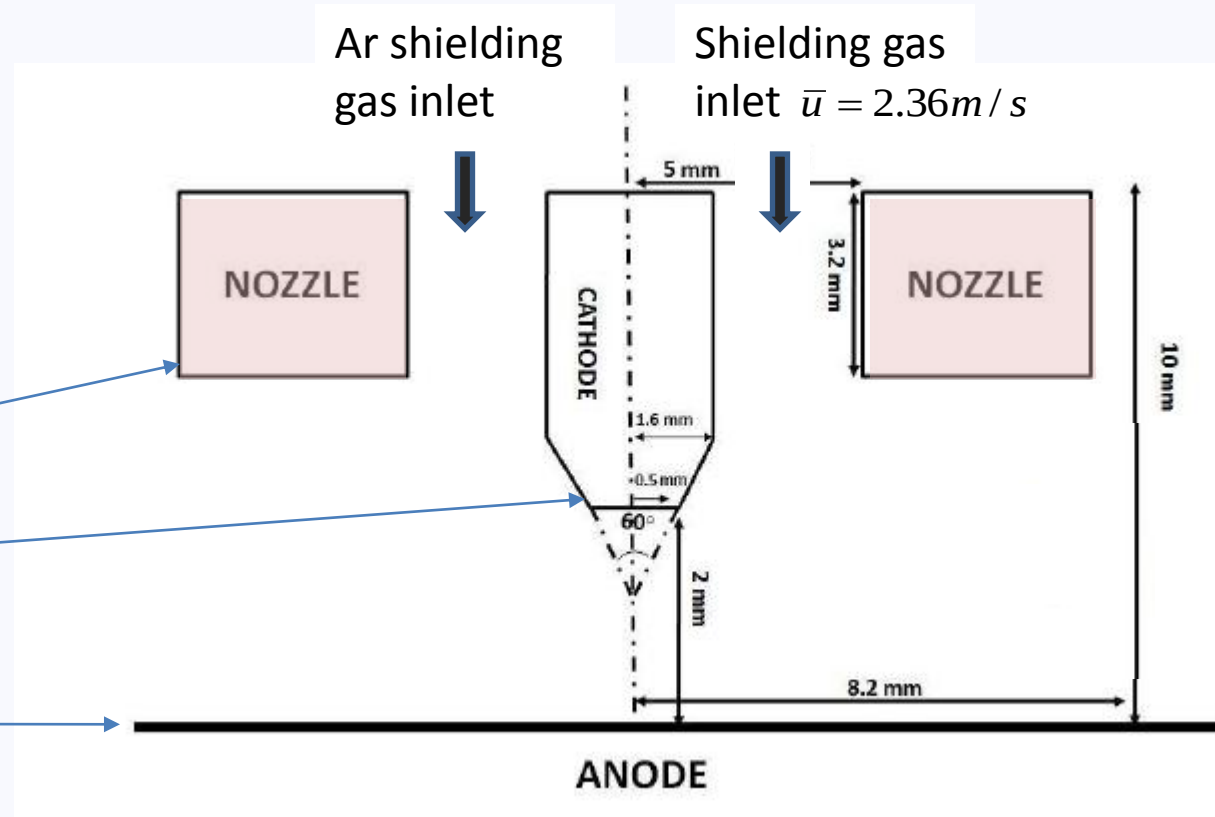
$\sigma = 2700 \text{ A}/(\text{V.m})$
 $\sigma = 10^{-5} \text{ A}/(\text{V.m})$



Azimuthal component of the magnetic field
along the radial direction ($r_0 = 10^{-3}$ m)

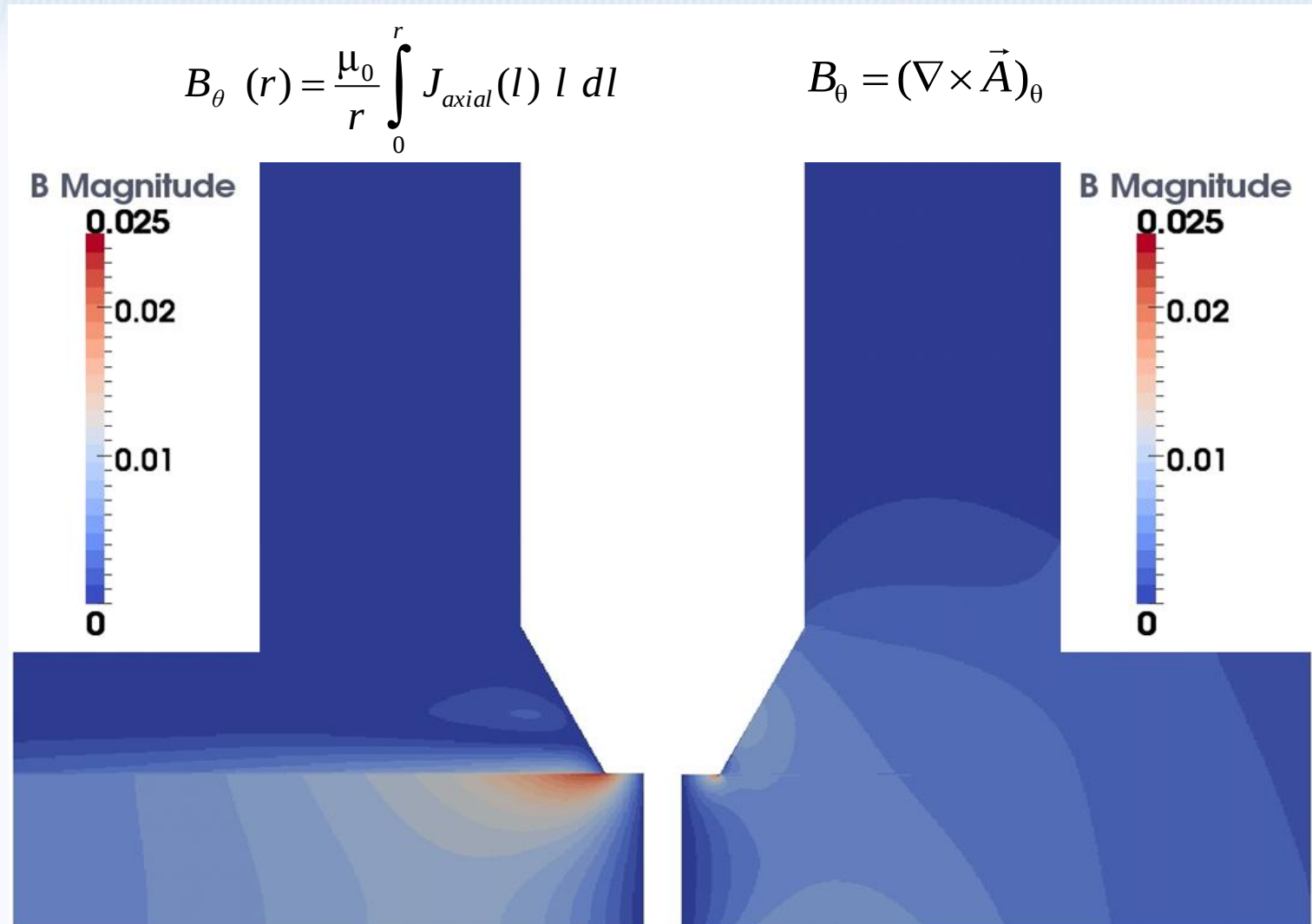
Test case: Tungsten Inert Gas welding

Applied current: $I=200\text{A}$



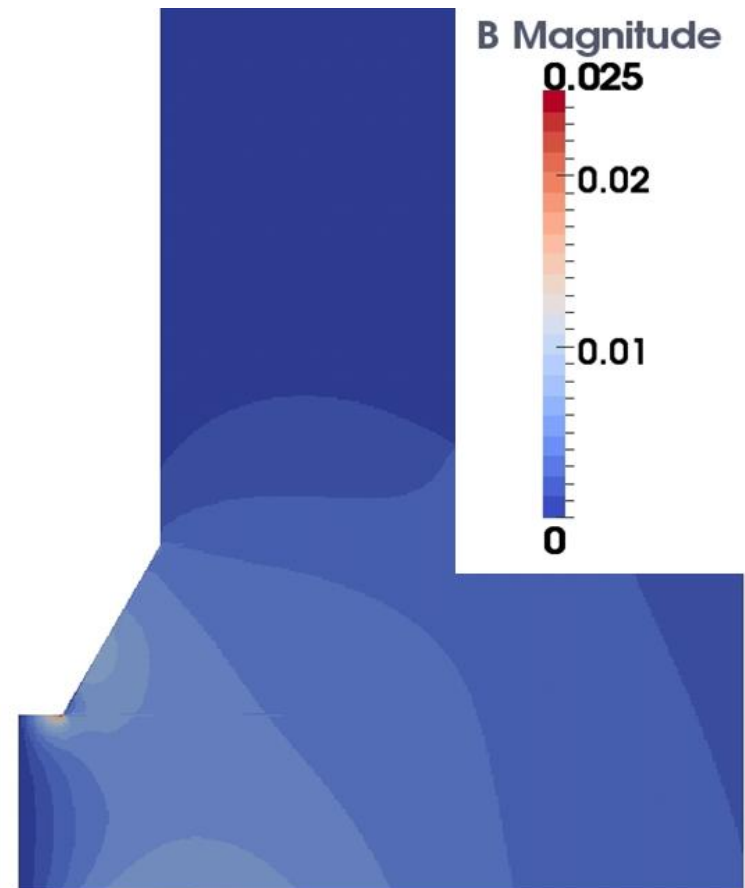
Picture of a TIG torch

Sketch of the cross section of a TIG torch



Magnetic field magnitude calculated with the electric potential formulation (left) and the axis-symmetric

$$B_{\theta} = (\nabla \times \vec{A})_{\theta}$$



Magnetic field magnitude
calculated with the axi-

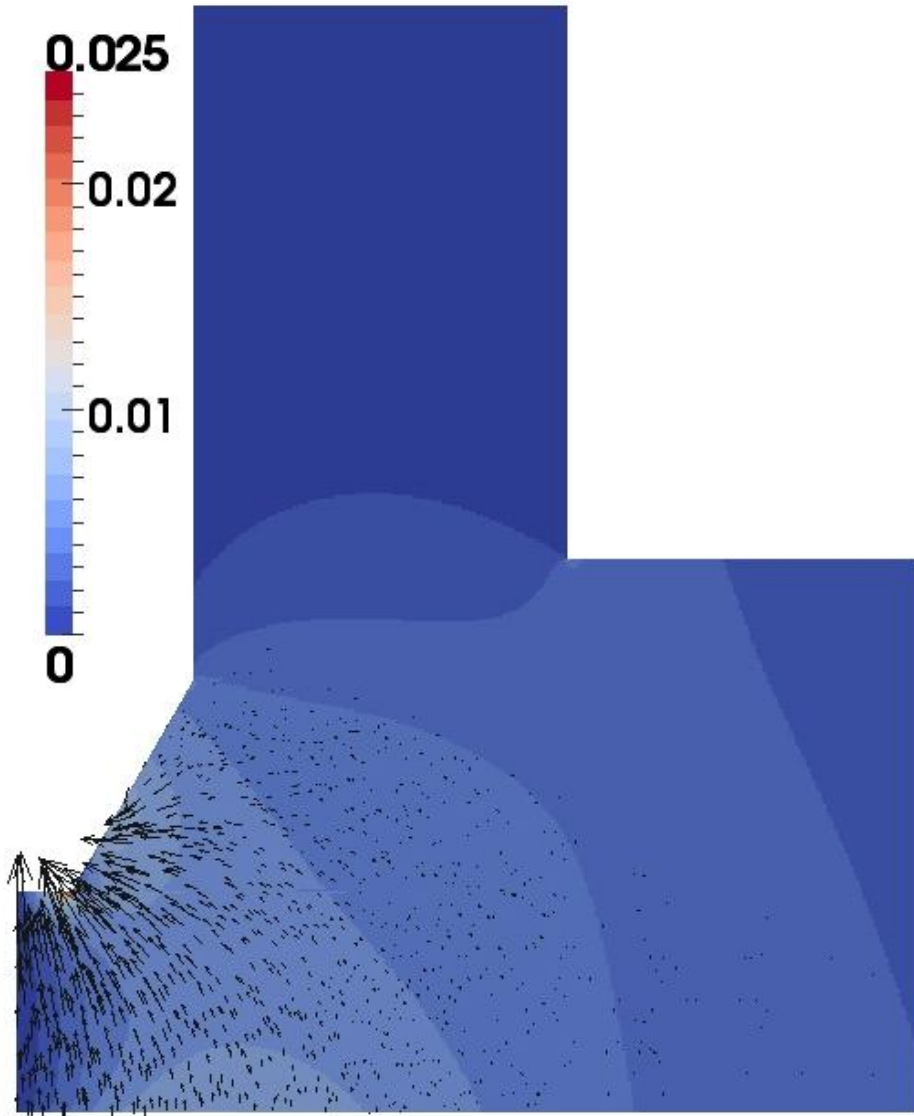
B Magnitude

0.025

0.02

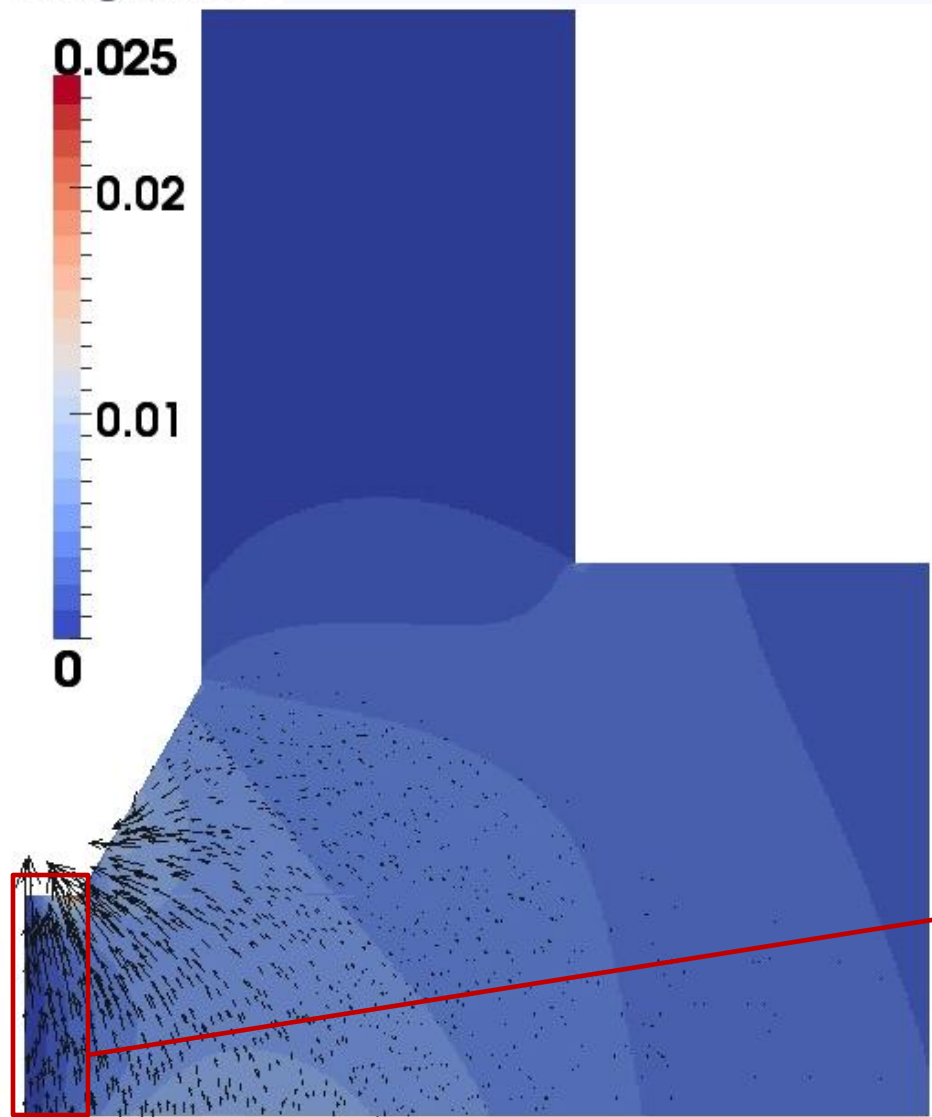
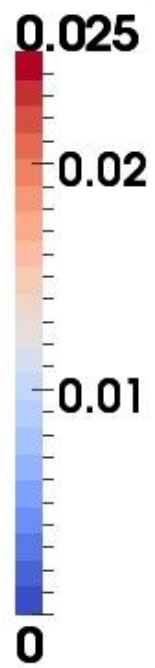
0.01

0

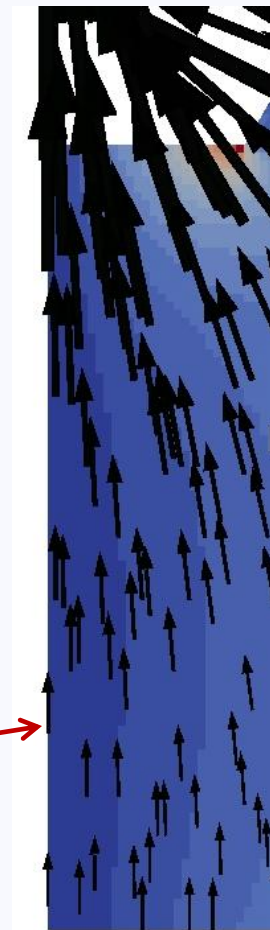


Current density $\vec{J}$

B Magnitude

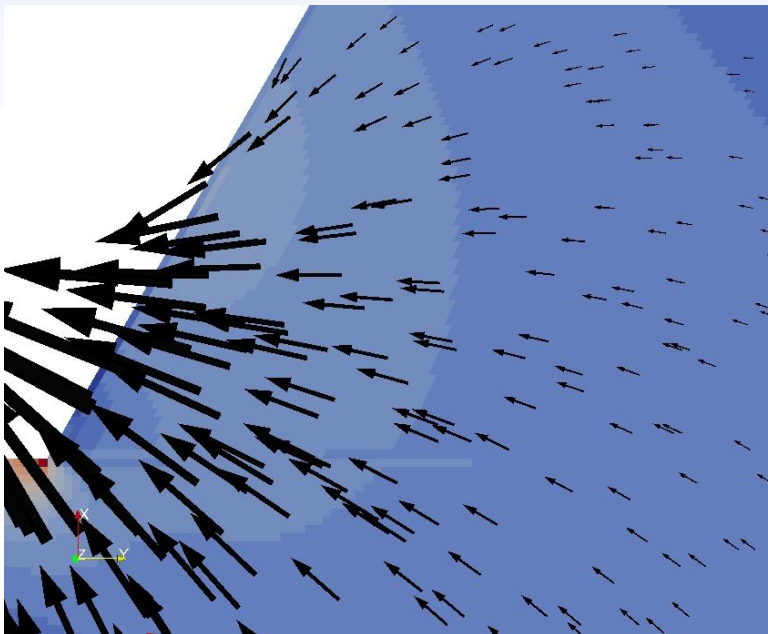
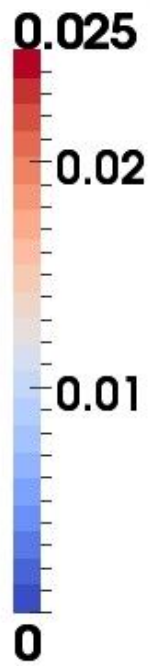


Current density $\vec{J}$



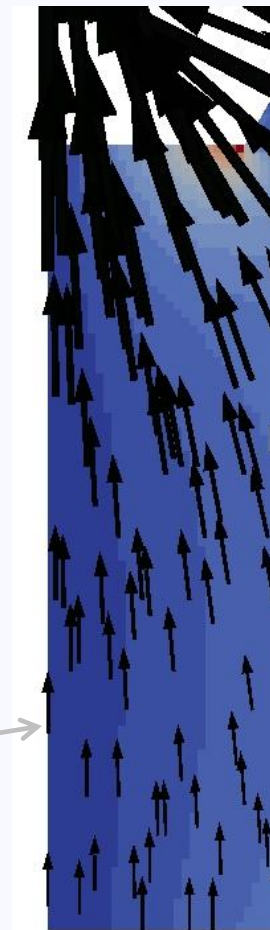
$$J_z \gg J_r$$

B Magnitude

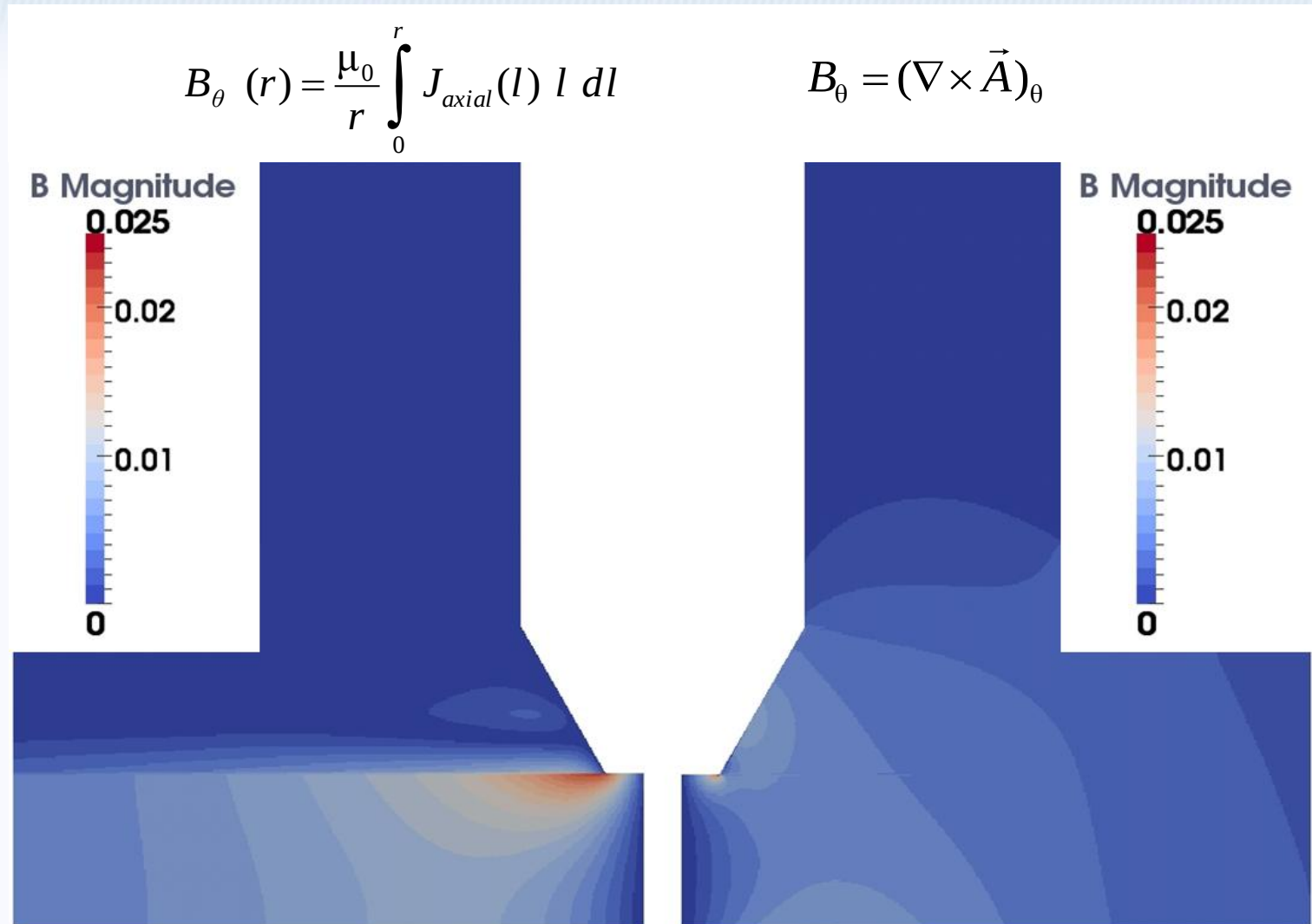


$$J_r \gg J_z$$

Current density $\vec{J}$



$$J_z \gg J_r$$



Magnetic field magnitude calculated with the electric potential formulation (left) and the axis-symmetric

electromagnetic model / conclusion

- 3D model with $V, \vec{A} \Rightarrow \vec{E}, \vec{J}, \vec{B}$
- 2D axi-symmetric models:
 - ✓ Electric potential formulation $V \Rightarrow \vec{E}, \vec{J} \Rightarrow B_\theta$
 - ✓ Magnetic field formulation $B_\theta \Rightarrow \vec{E}, \vec{J}$

Conclusions

- 3D model with $V, \vec{A} \rightarrow \vec{E}, \vec{J}, \vec{B}$

- 2D axi-symmetric models:



✓ Electric potential formulation if

✓ Magnetic field formulation has an "additional degree of freedom"

Conclusions

- 3D model with $V, \vec{A} \Rightarrow \vec{E}, \vec{J}, \vec{B}$

- 2D axi-symmetric models:

✓ $V, A_r, A_z \Rightarrow \vec{E}, \vec{J}, B_\theta$

✓

✓ Electric potential if

✓ Magnetic field formulation has an "additional degree of freedom"

Conclusions

- 3D model with $V, \vec{A} \Rightarrow \vec{E}, \vec{J}, \vec{B}$

- 2D axi-symmetric models:

✓ $V, A_r, A_z \Rightarrow \vec{E}, \vec{J}, B_\theta$

✓ $V, B_\theta \Rightarrow \vec{E}, \vec{J}$

✓ Electric potential formulation is simpler than vector potential if

✓ Magnetic field formulation has an "additional degree of freedom"

Conclusions

- 3D model with $V, \vec{A} \longrightarrow \vec{E}, \vec{J}, \vec{B}$

- 2D axi-symmetric models:

✓ $V, A_r, A_z \longrightarrow \vec{E}, \vec{J}, B_\theta$

✓ $V, B_\theta \longrightarrow \vec{E}, \vec{J}$

✓ Electric potential formulation $V \longrightarrow \vec{E}, \vec{J} \longrightarrow B_\theta$ if

✓ Magnetic field formulation has an "additional degree of freedom"

Conclusions

- 3D model with $V, \vec{A} \Rightarrow \vec{E}, \vec{J}, \vec{B}$

- 2D axi-symmetric models:

✓ $V, A_r, A_z \Rightarrow \vec{E}, \vec{J}, B_\theta$

✓ $V, B_\theta \Rightarrow \vec{E}, \vec{J}$

✓ Electric potential formulation $V \Rightarrow \vec{E}, \vec{J} \Rightarrow B_\theta$ if $J_r \ll J_z$

✓ Magnetic field formulation has an "additional degree of freedom"

Conclusions

- 3D model with $V, \vec{A} \Rightarrow \vec{E}, \vec{J}, \vec{B}$

- 2D axi-symmetric models:

✓ $V, A_r, A_z \Rightarrow \vec{E}, \vec{J}, B_\theta$

✓ $V, B_\theta \Rightarrow \vec{E}, \vec{J}$

✓ Electric potential formulation $V \Rightarrow \vec{E}, \vec{J} \Rightarrow B_\theta$ if $J_r \ll J_z$

✓ Magnetic field formulation has an "additional degree of freedom"

Conclusions

- 3D model with $V, \vec{A} \Rightarrow \vec{E}, \vec{J}, \vec{B}$

- 2D axi-symmetric models:

✓ $V, A_r, A_z \Rightarrow \vec{E}, \vec{J}, B_\theta$

✓ $V, B_\theta \Rightarrow \vec{E}, \vec{J}$

✓ Electric potential formulation $V \Rightarrow \vec{E}, \vec{J} \Rightarrow B_\theta$ if $J_r \ll J_z$

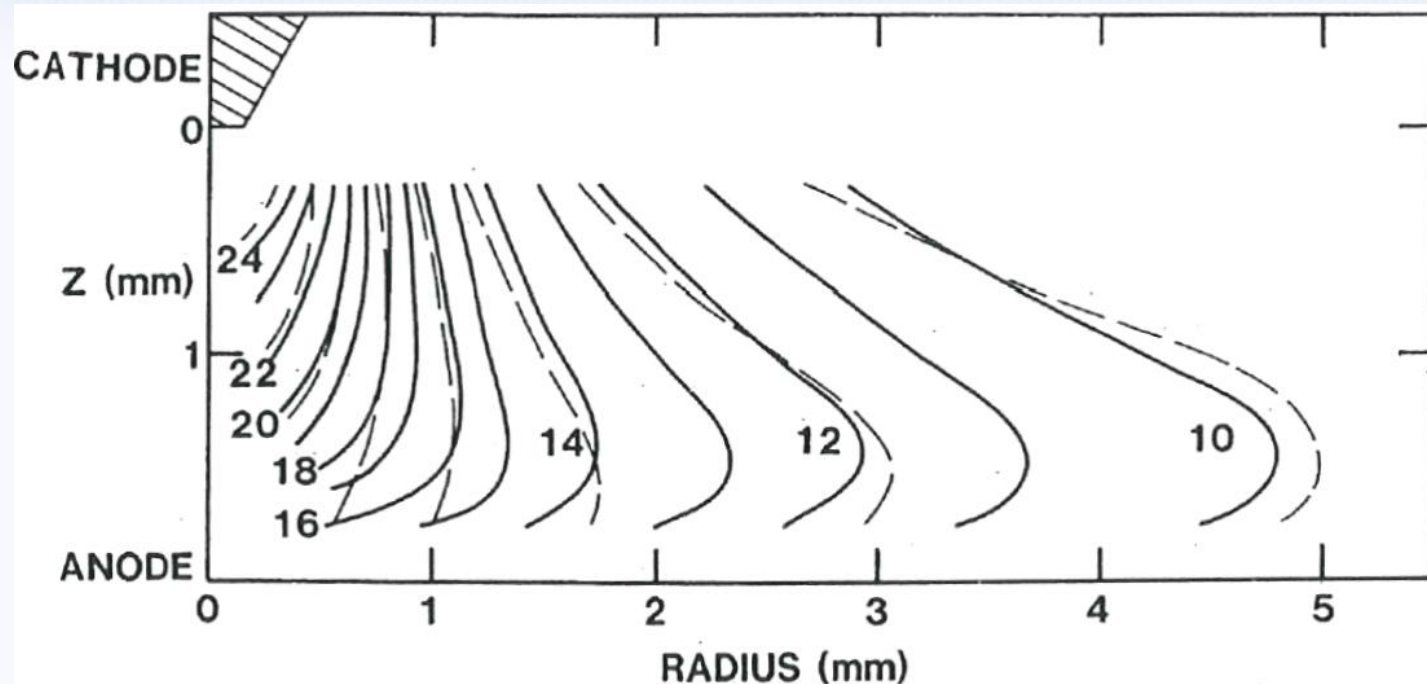
✓ Magnetic field formulation has an "additional degree of freedom"

- Acknowledgements

To Profs. Jacques Aubreton and Marie-Françoise Elchinger
for the data tables of thermodynamic and transport
properties.

To KK-foundation,
ESAB, and
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Thank you for your attention !



Measured temperature profile for
a current intensity 200 A and 2 mm long
arc

Figure from: *G.N. Haddad and A.J.D. Farmer (1985). Temperature measurements in gas tungsten arcs, Welding J, 64, pp. 339-342.*

Boundary conditions:

M.C. Tsai, and Sindo Kou (1990). Heat transfer and fluid flow in welding arcs produced by sharpened and flat electrodes, Int. J. Heat Mass Transfer, 33, pp. 2089-2098